Universal Spectral Behavior of $x^2(ix)^{\varepsilon}$ Potentials (work with Carl Bender)

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In 1998, Bender and Boettcher studied a class of Hamiltonians:

$$H = p^2 + x^2 (ix)^{\varepsilon}$$

where ε is taken as a real parameter.

The prescription to examine this Hamiltonian is well-known:

- convert into a differential equation form by replacing p by -id/dx

$$-\frac{d^2}{dx^2}\psi(x) + x^2(ix)^{\epsilon}\psi(x) = E\psi(x)$$

- Solve the eigenvalue problem numerically.

We are all familiar with the solution of the eigenvalue problem:

$$-\frac{d^2}{dx^2}\psi(x) + x^2(ix)^{\epsilon}\psi(x) = E\psi(x)$$





Isospectral Pairs of Hamiltonians

$$H_n = \eta p^n - \gamma (ix)^{n^2}$$

For example:

$$H_2 = \frac{1}{2m}p^2 - \gamma x^4 \qquad K_2 = \frac{1}{2m}x^2 + 4\gamma p^4 + \hbar \sqrt{\frac{2\gamma}{m}}p$$







Isospectral Pairs of Hamiltonians

- We've already seen 4th order Hamiltonians...
- We know some of the qualities of the classical paths
- So...can we construct a general class of fourth order oscillators that have a real eigenspectrum? (i.e. is PT general or is it special to second order problems?)

$$\left(\frac{d}{dx}\right)^4 \psi(x) = E\psi(x) - x^2(ix)^{\epsilon}\psi(x)$$
$$H = p^4 + x^2(ix)^{\epsilon}$$









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Four possible asymptotic behaviours:

$$\psi(x) \sim \exp\left(\frac{4\omega}{6+\varepsilon}x^{(6+\varepsilon)/4}e^{i\pi\varepsilon/8}\right) \quad (|x| \to \infty)$$

$$\omega^4 = -1$$

Stokes' wedges

Ignore "sign" so that the phase angle is: $(6 + \varepsilon)\theta/4 + \pi\varepsilon/8$ Centre of wedge: $-\pi\varepsilon/(12 + 2\varepsilon)$ Angular opening: $4\pi/(6 + \varepsilon)$

Worked Example

At $\varepsilon = 0$:

Opening angle is $2\pi/3$

As $\varepsilon \to \infty$:

Opening angle vanishes.

At $\varepsilon = -1$:

 $\theta = \pi/10$ and $\Delta\theta/2 = 4\pi/10$

Thus, the upper edge of the right wedge is at $\pi/2$.

Numerical Procedure

Substitute into differential equation:

$$\psi_R^{(4)}(r) = -\left(r^2 + Ee^{4i\theta_R}\right)\psi_R(r) = -Q(r)\psi_R(r)$$

Calculate WKB, so that we have initial conditions: $\psi_{\text{WKB}}(r_0) \sim KQ^{-\frac{3}{8}}(r_0) \exp\left[\omega \int_0^{r_0} dr \, Q^{1/4}(r)\right]$

Choose two solutions:

$$\omega_1 = e^{3\pi i/4} \quad \text{and} \quad \omega_2 = e^{-3\pi/4}$$

Numerical Procedure

Require continuity:

 $\begin{aligned} \alpha_1 \psi_{L,1}(0) + \alpha_2 \psi_{L,2}(0) &= \beta_1 \psi_{R,1}(0) + \beta_2 \psi_{R,2}(0), \\ \alpha_1 e^{-i\theta_L} \psi'_{L,1}(0) + \alpha_2 e^{-i\theta_L} \psi'_{L,2}(0) &= \beta_1 e^{-i\theta_R} \psi'_{R,1}(0) + \beta_2 e^{-i\theta_R} \psi'_{R,2}(0), \\ \alpha_1 e^{-2i\theta_L} \psi''_{L,1}(0) + \alpha_2 e^{-2i\theta_L} \psi''_{L,2}(0) &= \beta_1 e^{-2i\theta_R} \psi''_{R,1}(0) + \beta_2 e^{-2i\theta_R} \psi''_{R,2}(0), \\ \alpha_1 e^{-3i\theta_L} \psi'''_{L,1}(0) + \alpha_2 e^{-3i\theta_L} \psi'''_{L,2}(0) &= \beta_1 e^{-3i\theta_R} \psi'''_{R,1}(0) + \beta_2 e^{-3i\theta_R} \psi'''_{R,2}(0). \end{aligned}$

By Cramer's Rule, set Det = 0

 $\begin{pmatrix} \psi_{R,1}(0) & \psi_{R,2}(0) & -\psi_{L,1}(0) & -\psi_{L,2}(0) \\ e^{-i\theta_R}\psi'_{R,1}(0) & e^{-i\theta_R}\psi'_{R,2}(0) & -e^{-i\theta_L}\psi'_{L,1}(0) & -e^{-i\theta_L}\psi'_{L,2}(0) \\ e^{-2i\theta_R}\psi''_{R,1}(0) & e^{-2i\theta_R}\psi''_{R,2}(0) & -e^{-2i\theta_L}\psi''_{L,1}(0) & -e^{-2i\theta_L}\psi''_{L,2}(0) \\ e^{-3i\theta_R}\psi'''_{R,1}(0) & e^{-3i\theta_R}\psi'''_{R,2}(0) & -e^{-3i\theta_L}\psi'''_{L,1}(0) & -e^{-3i\theta_L}\psi'''_{L,2}(0) \end{pmatrix}$





















