

Non-Hermitian operators in quantum physics



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First, there's the room you can see through the glass – that's just the same as our drawing-room, only the things go the other way.

Hors de ligne

(Outline)

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1. QM with non-Hermitian operators
(just some conceptual remarks)
2. \mathcal{PT} -symmetry
(what is known and my point of view)
3. physical \mathcal{PT} -symmetric models in QM
(non-self-adjoint Robin boundary conditions)
4. imaginary cubic oscillator
(about the non-existence of the metric operator)
5. Conclusions

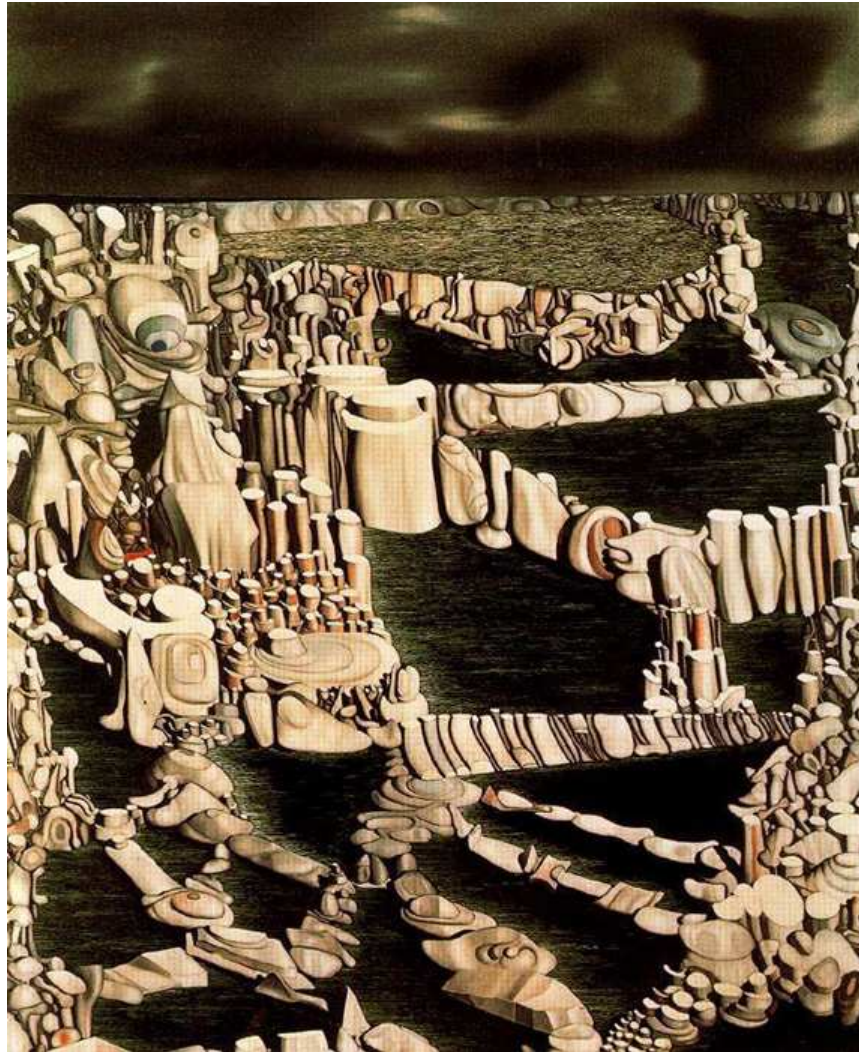


¿ QM with non-Hermitian operators ?

\mathbb{C}

\mathbb{R}

$$H^* = H$$



\mathbb{I}

$$H^{\mathcal{PT}} = H$$

Imaginary Numbers by Yves Tanguy, 1954
(Museo Thyssen-Bornemisza, Madrid)

Insignificant non-Hermiticity

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Theorem (spectral theorem).

Let $H = H^*$. Then

$$f(H) = \int_{\sigma(H)} f(\lambda) \, dE_H(\lambda)$$

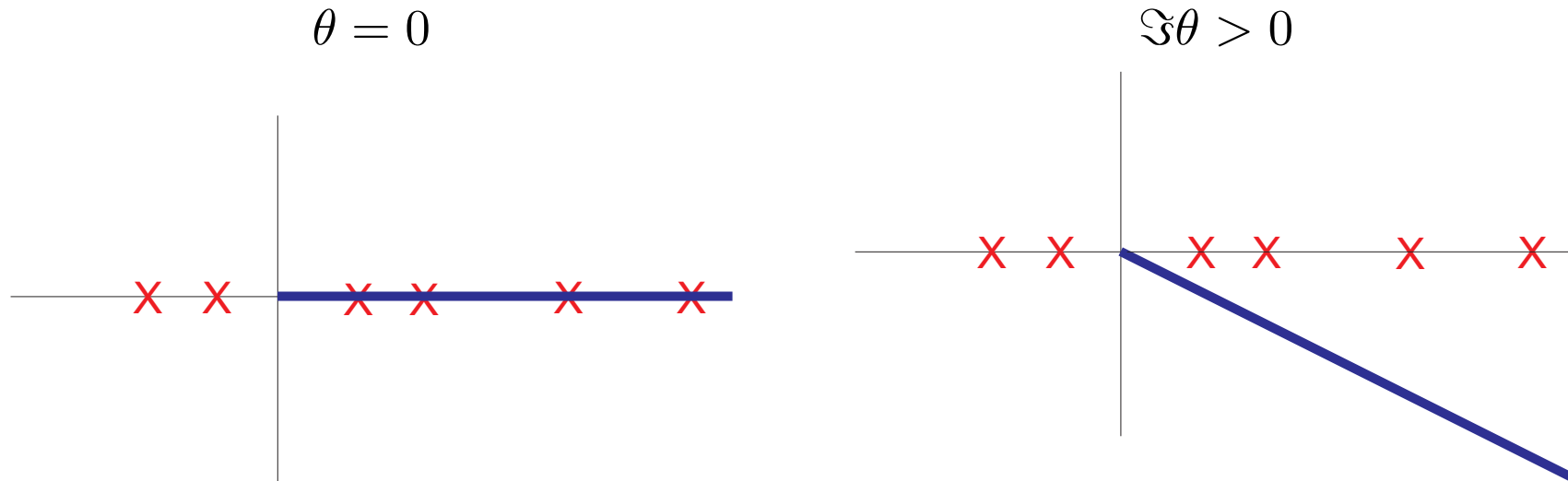
for any complex-valued continuous function f .



Technical non-Hermiticity

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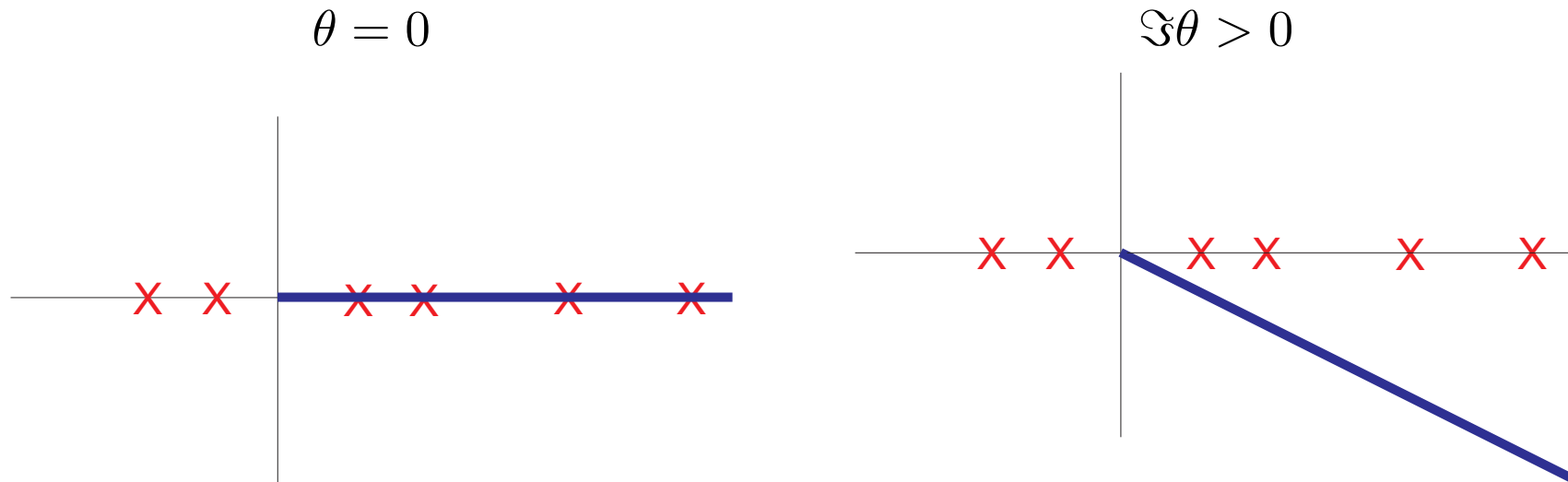
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[Aguilar/Balslev, Combes 1971], [Simon 1972], [Van Winter 1974], ...

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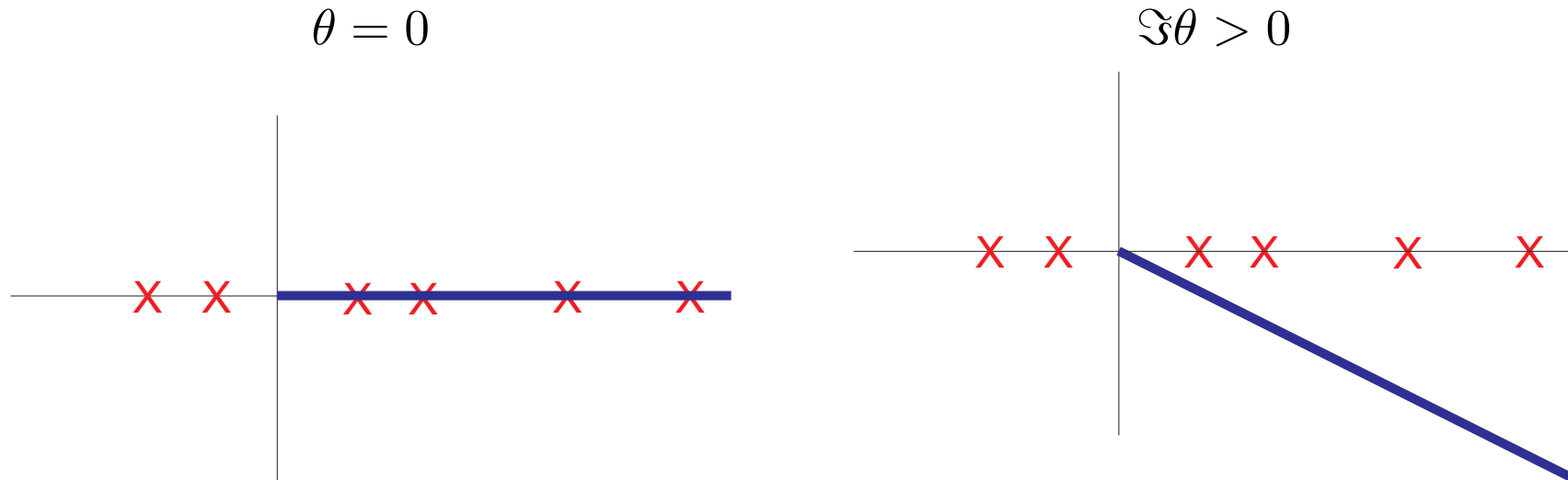
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Example 2. adiabatic transition probability for $H(t) := \vec{\gamma}(t/\tau) \cdot \vec{\sigma}$, $\tau \rightarrow \infty$

[Berry 1990], [Joye, Kunz, Pfister 1991], [Jakšić, Segert 1993], ...

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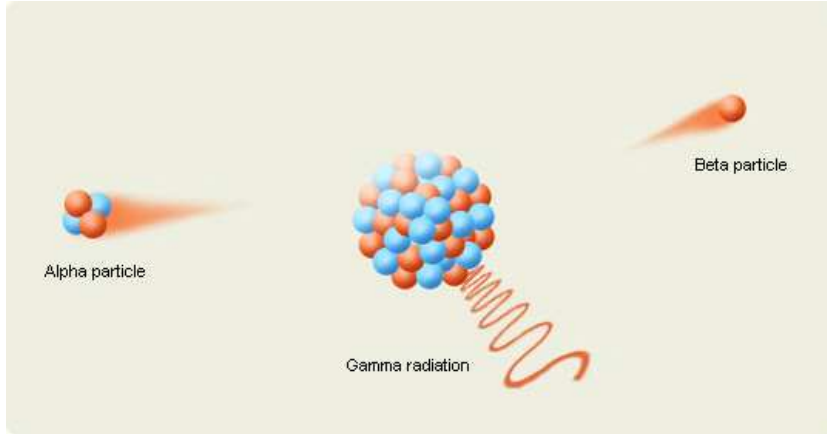
Example 3. Regge theory $H_l := -\frac{d^2}{dr^2} + \frac{l(l+1)}{r^2} + V(r)$, $l \in \mathbb{C}$

[Regge 1957], [Connor 1990], [Sokolovski 2011], ...

Approximate non-Hermiticity

open systems

Example 1. radioactive decay



Example 2. dissipative Schrödinger operators in semiconductor physics

Baro, Behrndt, Kaiser, Neidhardt, Rehberg, ...

Example 3. repeated interaction quantum systems

Bruneau, Joye, Merkli, Pillet, ...

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i.e. non-Hermitian **observables**,
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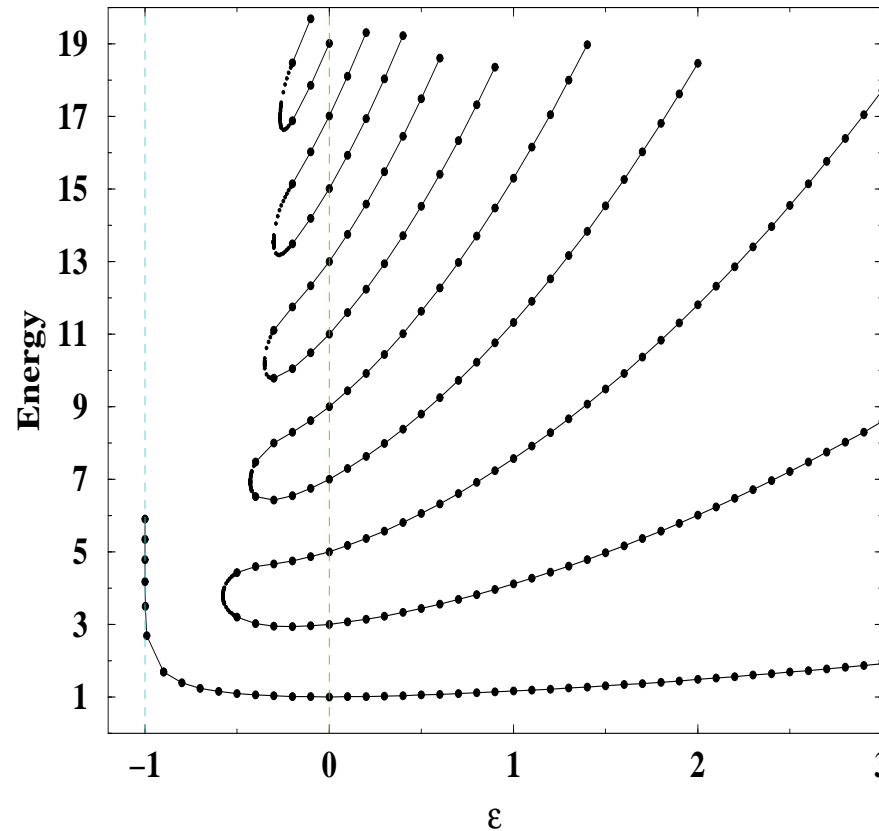
by changing the Hilbert space,
preserving a **similarity** to self-adjoint operators

Non-Hermitian Hamiltonians with real spectra

$$-\Delta + V \quad \text{in} \quad L^2(\mathbb{R})$$

$$V(x) = x^2 + ix^3$$

[Caliceti, Graffi, Maioli 1980]



$$V(x) = x^2 (ix)^\varepsilon$$

[Bessis, Zinn-Justin]

[Bender, Boettcher 1998]

[Dorey, Dunning, Tateo 2001]

[Shin 2002]

[Azizov, Kuzhel, Günther, Trunk 2010]

$$V(x) = \begin{cases} i \operatorname{sgn}(x) & \text{if } x \in (-L, L) \\ \infty & \text{elsewhere} \end{cases}$$

[Znojil 2001]

¿ What is behind the reality of the spectrum ?



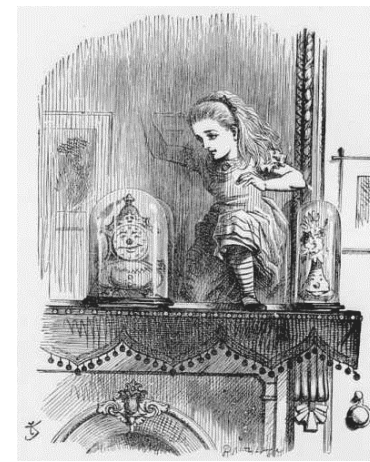
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$$[H, \mathcal{PT}] = 0$$

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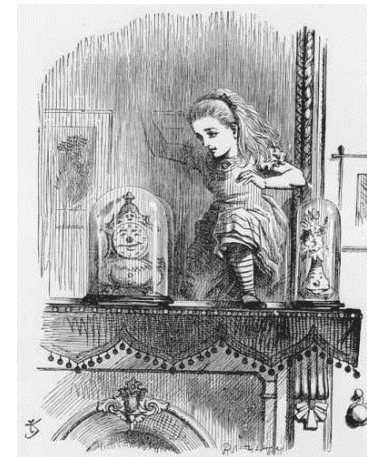


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$\implies H$ is **self-adjoint** in $(L^2, \langle \cdot, \Theta \cdot \rangle)$,
metric i.e. $\Theta^{1/2} H \Theta^{-1/2}$ is self-adjoint in $(L^2, \langle \cdot, \cdot \rangle)$



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metric

Albeverio-Fei-Kurasov, Bender-Brody-Jones, Caliceti-Graffi-Sjöstrand, Fring, Graefe-Schubert, Kretschmer-Szymanowski, Langer-Tretter, Mostafazadeh, Scholtz-Geyer-Hahne, Znojil, ...

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concept of “unbounded metric” is trivial and physically doubtful

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Remark. In general (in ∞ -dimensional spaces), all the classes of operators are **unrelated**.

[Siegl 2008]

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suggestions:

- nuclear physics [Scholtz, Geyer, Hahne 1992]
- optics [Klaيمان, Günther, Moiseyev 2008], [Schomerus 2010], [West, Kottos, Prosen 2010]
- solid state physics [Bendix, Fleischmann, Kottos, Shapiro 2009]
- superconductivity [Rubinstein, Sternberg, Ma 2007]
- electromagnetism [Ruschhaupt, Delgado, Muga 2005], [Mostafazadeh 2009]

experiments:

- optics [Guo *et al.* 2009], [Longhi 2009], [Rüter *et al.* 2010]
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¡ but !

*“So far, there have been no experiments that prove clearly and definitively that **quantum** systems defined by non-Hermitian \mathcal{PT} -symmetric Hamiltonians do exist in nature.”*

[Bender 2007]

The simplest \mathcal{PT} -symmetric model

[D.K., Bíla, Znojil 2006]

$$\mathcal{H} := L^2\left(-\frac{\pi}{2}, \frac{\pi}{2}\right)$$

$$H_\alpha \psi := -\psi'', \quad D(H_\alpha) := \left\{ \psi \in W^{2,2}\left(-\frac{\pi}{2}, \frac{\pi}{2}\right) \mid \psi'(\pm \frac{\pi}{2}) + i\alpha \psi(\pm \frac{\pi}{2}) = 0 \right\}, \quad \alpha \in \mathbb{R}$$

$-\Delta$


$\frac{d\psi}{dn} - i\alpha\psi = 0$ $\frac{d\psi}{dn} + i\alpha\psi = 0$

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A horizontal line segment with solid black square endpoints. Above the segment is the label $-\Delta$. Below the left endpoint is the boundary condition $\frac{d\psi}{dn} - i\alpha\psi = 0$. Below the right endpoint is the boundary condition $\frac{d\psi}{dn} + i\alpha\psi = 0$.

Theorem 1. H_α is an m -sectorial operator with compact resolvent satisfying

$$H_\alpha^* = H_{-\alpha} = \mathcal{T}H_\alpha\mathcal{T} \quad (\mathcal{T}\text{-self-adjointness})$$

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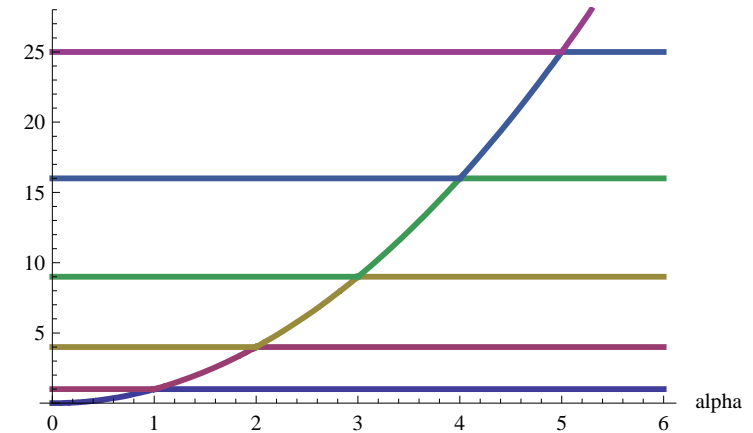
$$\begin{array}{c} \text{---} \overline{\Delta} \text{---} \\ \frac{d\psi}{dn} - i\alpha\psi = 0 \qquad \frac{d\psi}{dn} + i\alpha\psi = 0 \end{array}$$

Theorem 1. H_α is an m -sectorial operator with compact resolvent satisfying

$$H_\alpha^* = H_{-\alpha} = \mathcal{T}H_\alpha\mathcal{T} \quad (\mathcal{T}\text{-self-adjointness})$$

Theorem 2.

$$\sigma(H_\alpha) = \{\alpha^2\} \cup \{n^2\}_{n=1}^\infty$$



The simplest \mathcal{PT} -symmetric model

[D.K., Bíla, Znojil 2006]

$$\mathcal{H} := L^2(-\frac{\pi}{2}, \frac{\pi}{2})$$

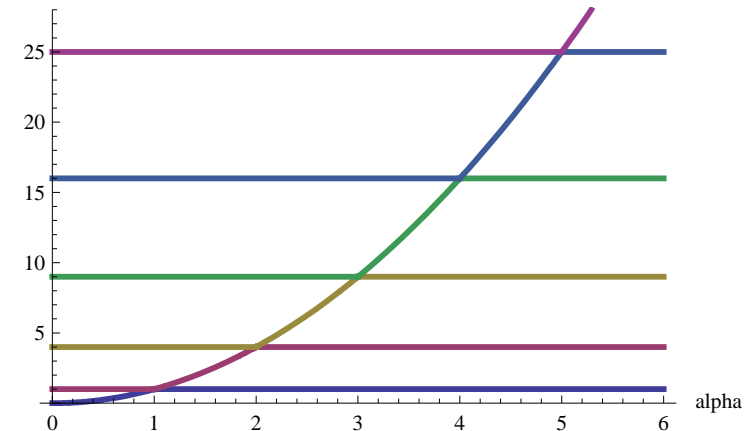
$$H_\alpha \psi := -\psi'', \quad D(H_\alpha) := \left\{ \psi \in W^{2,2}(-\frac{\pi}{2}, \frac{\pi}{2}) \mid \psi'(\pm \frac{\pi}{2}) + i\alpha \psi(\pm \frac{\pi}{2}) = 0 \right\}, \quad \alpha \in \mathbb{R}$$

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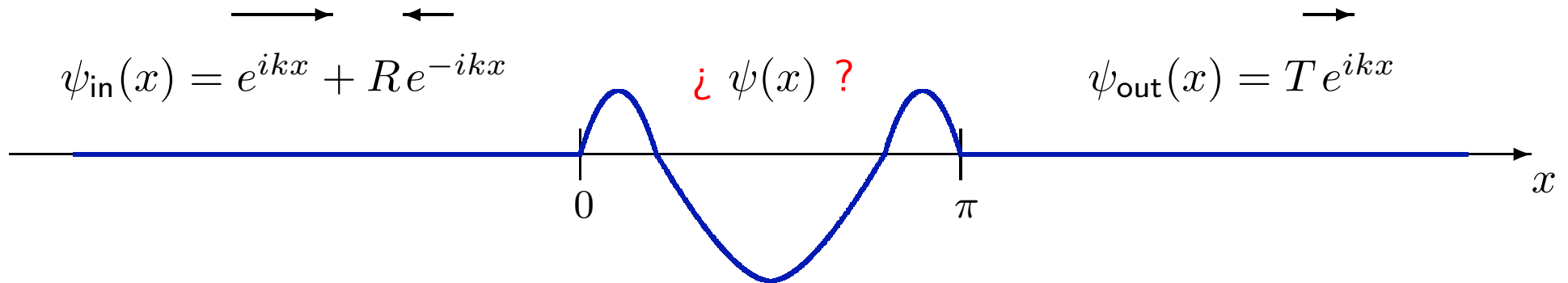


Corollary. The spectrum of H_α is $\begin{cases} \text{always real,} \\ \text{simple if } \alpha \notin \mathbb{Z} \setminus \{0\}. \end{cases}$

Scattering realisation in QM

[Hernandez-Coronado, D.K., Siegl 2011]

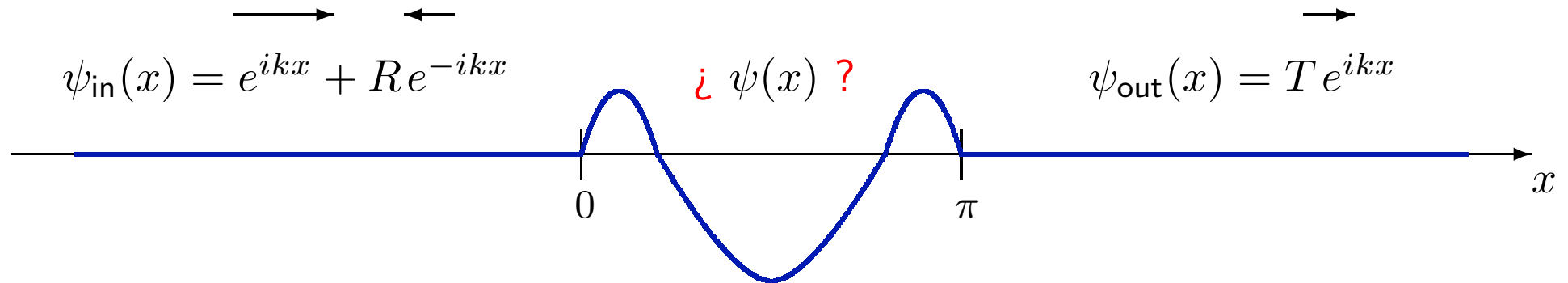
scattering by a compactly supported even potential V : $-\psi'' + V\psi = k^2\psi \quad k > 0$



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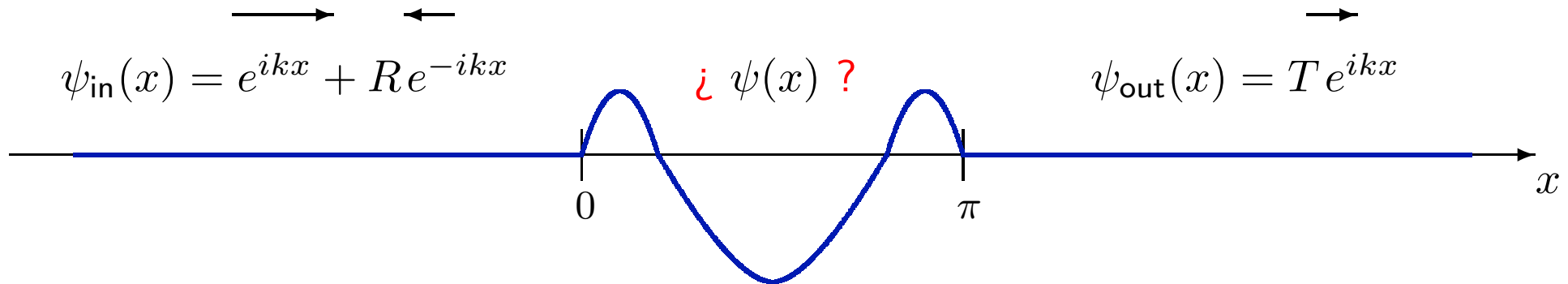


perfect transmission $\implies \begin{cases} -\psi'' + V\psi = k^2\psi & \text{in } (0, \pi) \\ \psi' - ik\psi = 0 & \text{at } 0, \pi \end{cases}$ non-linear
(i.e. $R = 0$)

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solutions given by a non-self-adjoint \mathcal{PT} -symmetric spectral problem:

$$\begin{cases} -\psi'' + V\psi = \mu(\alpha)\psi & \text{in } (0, \pi) \\ \psi' + i\alpha\psi = 0 & \text{at } 0, \pi \\ \mu(\alpha) = \alpha^2 \end{cases}$$

non-Hermitian



The metric operator

[D.K., Siegl, Železný 2011]

Theorem 3. Let $\alpha \notin \mathbb{Z} \setminus \{0\}$.

Then H_α is similar to a self-adjoint operator $h_\alpha := \Omega H_\alpha \Omega^{-1}$ with the metric

$$\Theta := \Omega^* \Omega = I + K$$

$$K(x, y) = \frac{2i}{\pi} e^{i\frac{\alpha}{2}(x-y)} \sin\left(\frac{\alpha}{2}(x-y)\right) + \frac{i\alpha}{\pi} (|y-x| - \pi) \operatorname{sgn}(y-x) + \frac{\alpha^2}{\pi} \left(\frac{\pi^2}{4} - xy - \frac{\pi}{2}|y-x|\right)$$

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Moreover, $h_\alpha = -\Delta_N + \alpha^2 \chi_0^N \langle \chi_0^N, \cdot \rangle$!!! $\chi_0^N(x) := \pi^{-1/2}$ (Neumann ground state)

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Proof. “backward usage of the spectral theorem” [D.K. 2008]

$$E_n = \begin{cases} \alpha^2 & \text{if } n = 0 \\ n^2 = E_n^N = E_n^D & \text{if } n \geq 1 \end{cases} \quad \phi_n(x) = \begin{cases} \frac{1}{\sqrt{\pi}} e^{i\alpha(x+\frac{\pi}{2})} & \text{if } n = 0 \\ \chi_n^N(x) + i \frac{\alpha}{n} \chi_n^D(x) & \text{if } n \geq 1 \end{cases}$$

$$\begin{aligned} \Omega &:= \sum_{n=0}^{\infty} \chi_n^N \langle \phi_n, \cdot \rangle = \chi_0^N \langle \phi_0, \cdot \rangle - \chi_0^N \langle \chi_0^N, \cdot \rangle + \sum_{n=0}^{\infty} \chi_n^N \langle \chi_n^N, \cdot \rangle - i\alpha \sum_{n=1}^{\infty} \frac{1}{n} \chi_n^N \langle \chi_n^D, \cdot \rangle \\ &= \chi_0^N \langle \phi_0, \cdot \rangle - \chi_0^N \langle \chi_0^N, \cdot \rangle + I + \alpha p (-\Delta_D)^{-1} \quad (-\Delta_D = p^* p) \end{aligned} \quad \text{q.e.d.}$$

The \mathcal{C} operator

Definition ([Bender, Brody Jones 2002], [Albeverio, Kuzhel 2005]).

Let H be \mathcal{P} -self-adjoint and \mathcal{C} bounded.

$$H \text{ is } \mathcal{C}\text{-symmetric} :\Longleftrightarrow \begin{cases} [H, \mathcal{C}] = 0 \\ \mathcal{C}^2 = I \\ \mathcal{P}\mathcal{C} \text{ is a metric for } H \end{cases}$$

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Our explicit result:

$\mathcal{C} = \mathcal{P} + L$ with

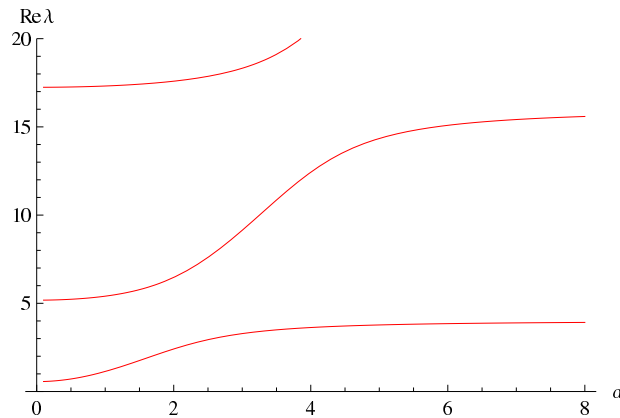
$$L(x, y) = \alpha e^{-i\alpha(y+x)} \left[\tan(\alpha \frac{\pi}{2}) - i \operatorname{sgn}(y+x) \right] \quad (|\alpha| < 1)$$

General \mathcal{PT} - symmetric case

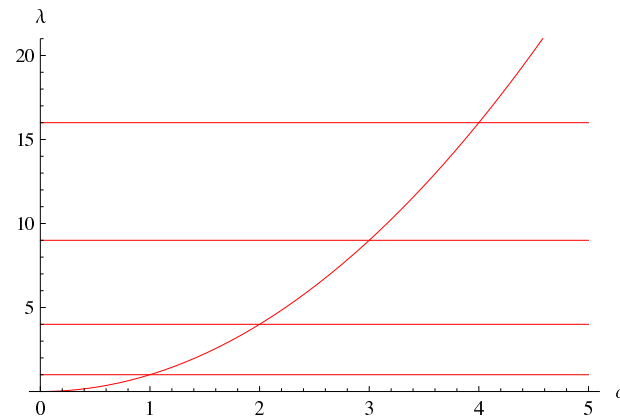
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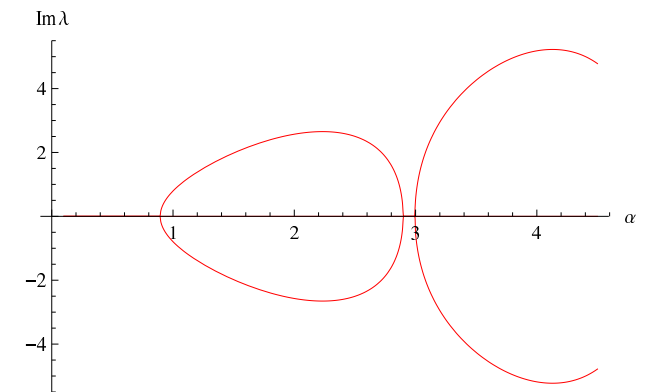
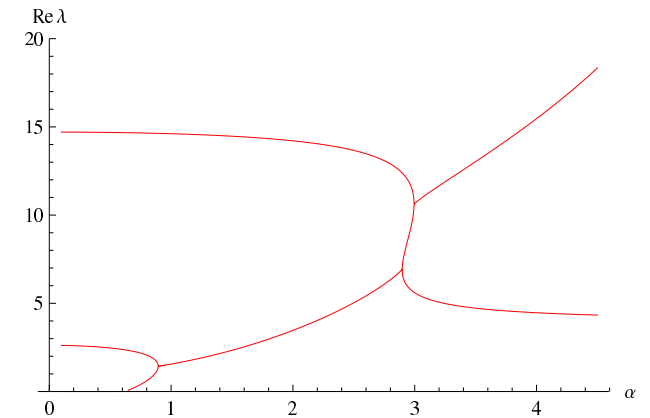
$\beta > 0$



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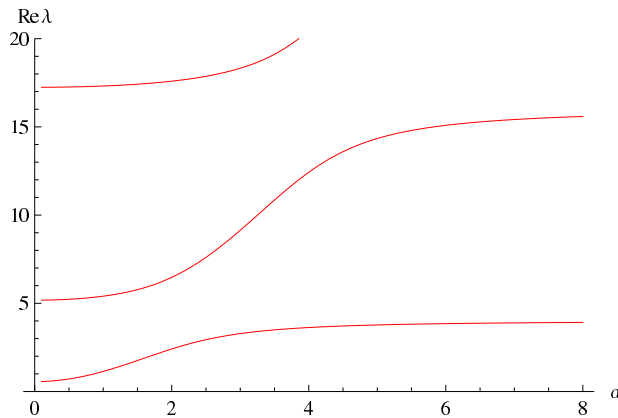


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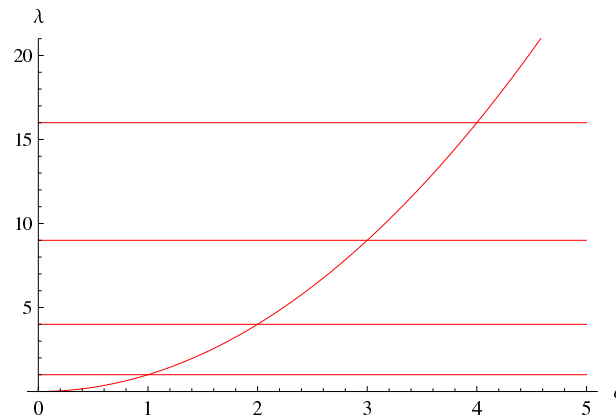
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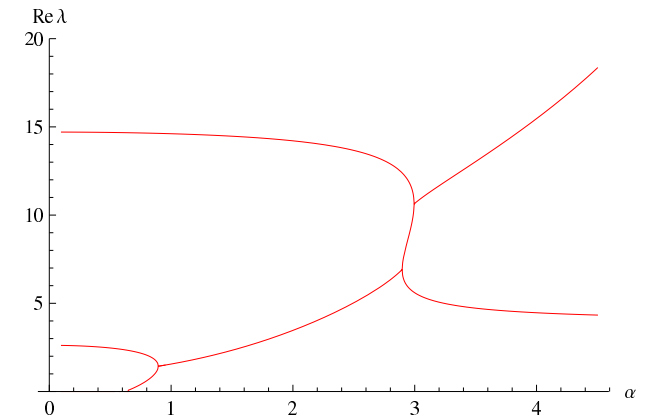
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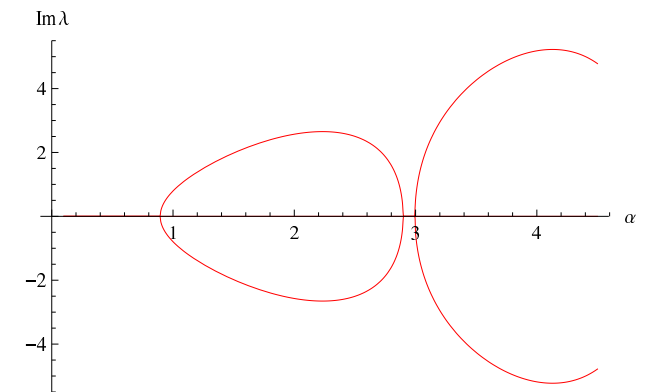
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$\Theta = I + K$ with

$$K(x, y) = e^{i\alpha(x-y) - \beta|x-y|} [c + i\alpha \operatorname{sgn}(x-y)] \quad c \in \mathbb{R}$$

$\Theta > 0$ if $\beta > 0$ large or $c^2 + \alpha^2$ small



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[D.K., Siegl 2012]

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- $\|(H_h - z)^{-1}\| = \mathcal{O}(h^{-n}), \quad \forall n > 0 \Rightarrow$ contradiction *q.e.d.*
[Davies 1999]

Pseudospectra and \mathcal{PT} -symmetry

[work in progress with Siegl and Tater]

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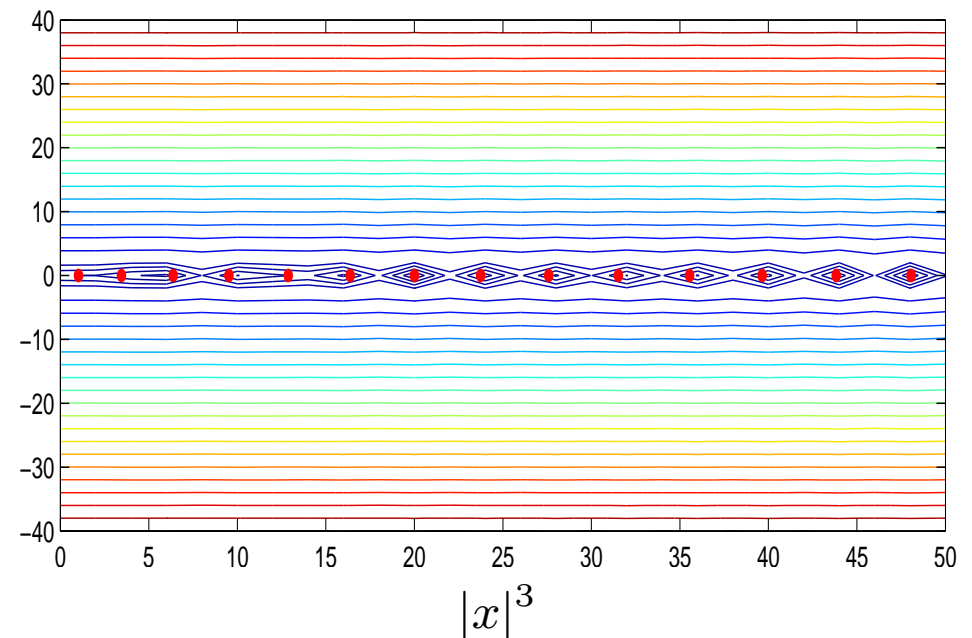
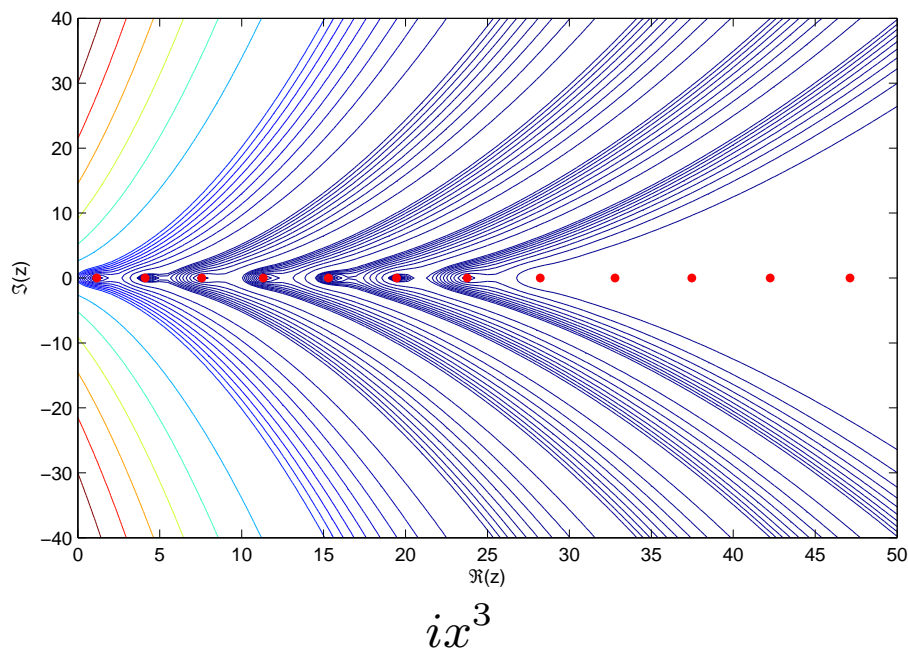
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- \exists metric \implies trivial pseudospectrum



Conclusions

Ad \mathcal{PT} -symmetry:

- no extension of QM
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Ad ix^3 :

- ! metric operator does not exist !

(bad basicity properties, non-trivial pseudospectrum, spectral instabilities)



Collection of open problems

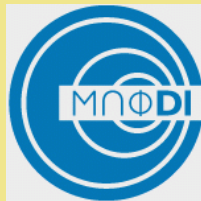
ESF exploratory workshop on

Mathematical aspects of the physics with non-self-adjoint operators

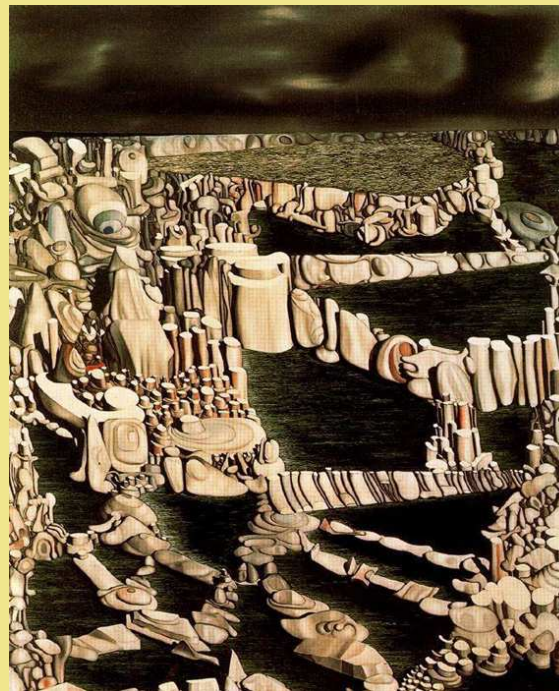
30 August - 3 September 2010

Prague, Czech Republic

<http://www.ujf.cas.cz/ESFxNSA/>



Doppler Institute



Imaginary numbers 1954 by Y. Tanguy

Studying non-self-adjoint operators is like being a vet rather than a doctor: one has to acquire a much wider range of knowledge, and to accept that one cannot expect to have as high a rate of success when confronted with particular cases.

E. B. Davies 2007

Convenors: J.-P. Gazeau (Paris), D. Krejcirik (Prague), P. Siegl (Prague)

Some of the open problems also available in *Integral Equations Operator Theory*.

My \mathcal{PT} -symmetric life

<http://gemma.ujf.cas.cz/~david/>

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- D. Kochan, D.K., R. Novák, P. Siegl: *The Pauli equation with complex boundary conditions*; J. Phys. A: Math. Theor., to appear; arXiv:1203.5011 [quant-ph] (2012).
- P. Siegl, D.K.: *Metric operator for the imaginary cubic oscillator does not exist*; arXiv:1208.1866 [quant-ph] (2012).