

Breakdown of adiabatic transfer schemes in the presence of decay or absorption

Eva-Maria Graefe

Department of Mathematics, Imperial College London, UK

joint work with

Nimrod Moiseyev and Alexei Mailybaev

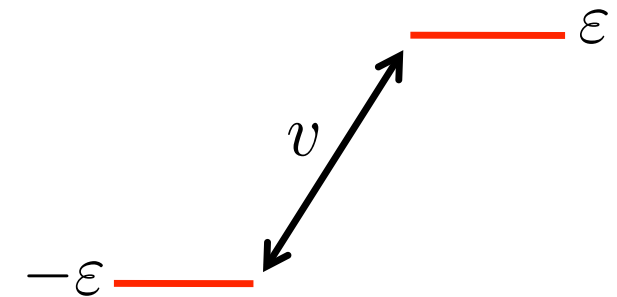
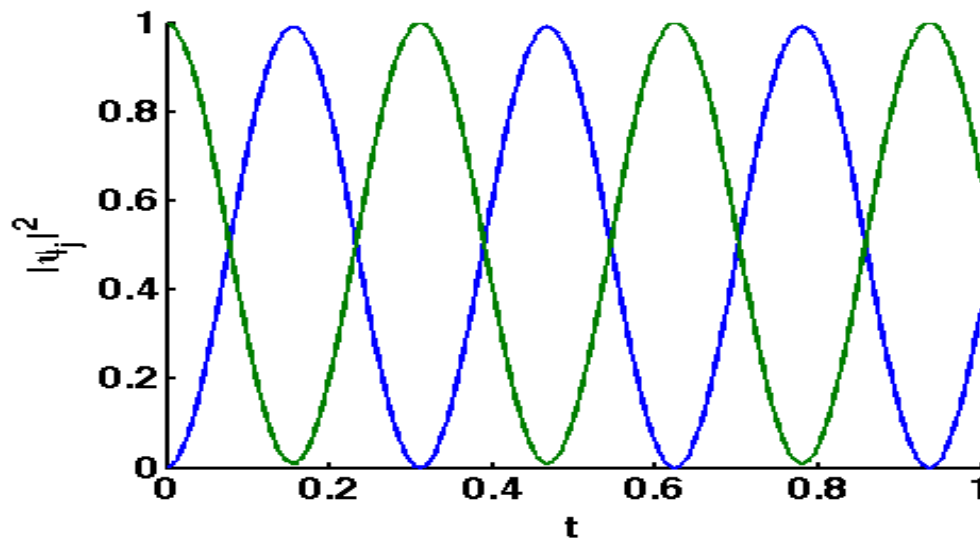
Technion, Haifa, Isreal

Instituto Nacional de Matemática Pura e Aplicada, Rio de Janeiro, Brazil

Quantum Population transfer

★ Task: Population transfer between two quantum states (e.g. in atomic system)

★ Coupling of the two states (e.g. via laser field of appropriate frequency)



$$H = \begin{pmatrix} \varepsilon & v \\ v & -\varepsilon \end{pmatrix}$$

★ Adiabatic state transfer: Slower, but efficient and robust

Adiabatic quantum evolution

Adiabatic theorem: System prepared in eigenstate stays in corresponding instantaneous eigenstate for infinitely slow parameter variation

★ Adiabatic state transfer: Slow parameter variation such that initial and target states are connected via instantaneous states

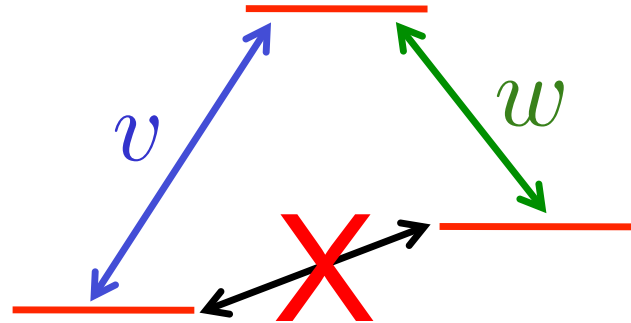
Adiabatic population transfer in three level systems - STIRAP

- ★ Population transfer in atomic systems between two states without direct transition



STIRAP

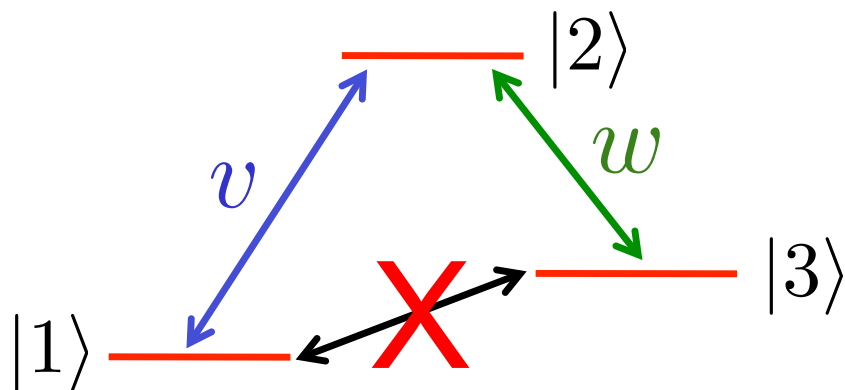
★ Population transfer in atomic systems between two states without direct transition



★ Via additional state

STIRAP

★ Population transfer in atomic systems between two states without direct transition

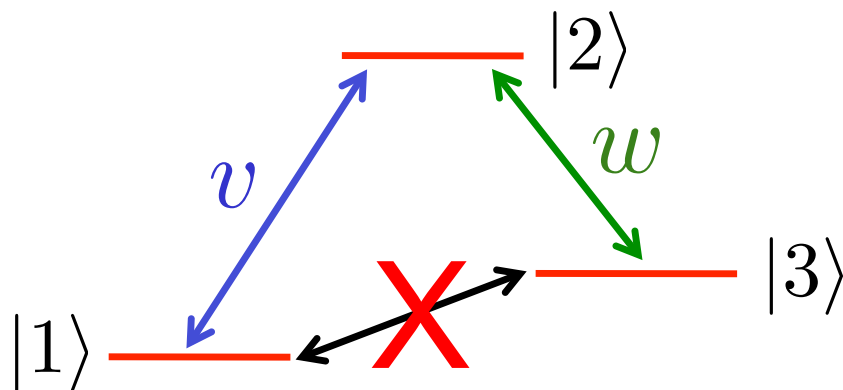


$$H = \begin{pmatrix} E_1 & v(t) & 0 \\ v(t) & E_2 & w(t) \\ 0 & w(t) & E_3 \end{pmatrix}$$

★ Via additional state

STIRAP

★ Population transfer in atomic systems between two states without direct transition



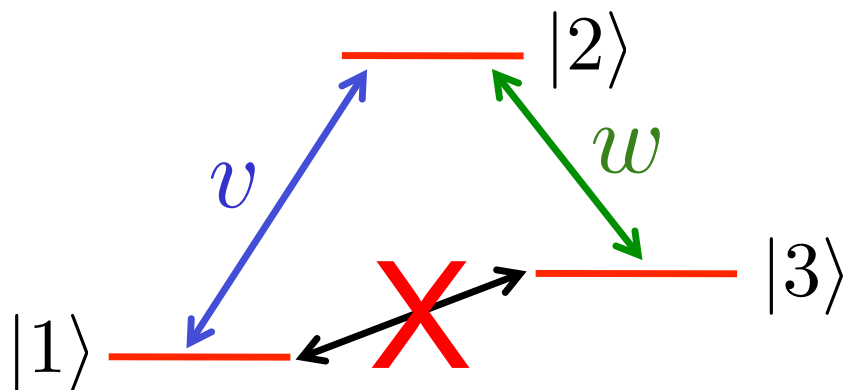
$$H = \begin{pmatrix} E_1 & v(t) & 0 \\ v(t) & E_2 & w(t) \\ 0 & w(t) & E_3 \end{pmatrix}$$

★ Via additional state

★ Intuitive scheme: First switch on v , then w . Nearly total transfer via state $|2\rangle$ possible (precise parameter control required)

STIRAP

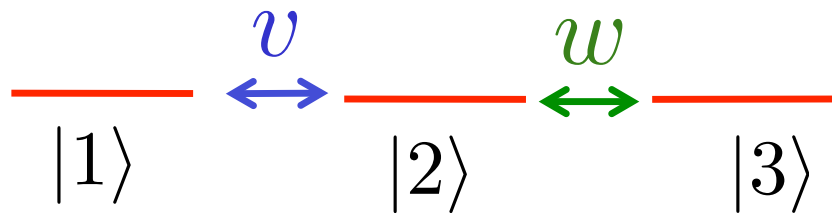
- ★ Population transfer in atomic systems between two states without direct transition



$$H = \begin{pmatrix} E_1 & v(t) & 0 \\ v(t) & E_2 & w(t) \\ 0 & w(t) & E_3 \end{pmatrix}$$

- ★ Via additional state
- ★ STIRAP: adiabatic nearly total population transfer from state $|1\rangle$ to state $|3\rangle$ without significant population of state $|2\rangle$ at any time

Eigenvalues and eigenstates

$$H(t) = \begin{pmatrix} 0 & v(t) & 0 \\ v(t) & 0 & w(t) \\ 0 & w(t) & 0 \end{pmatrix}$$


$|1\rangle \quad |2\rangle \quad |3\rangle$

$E_0 = 0$

$E_{\pm} = \pm \sqrt{v^2 + w^2}$

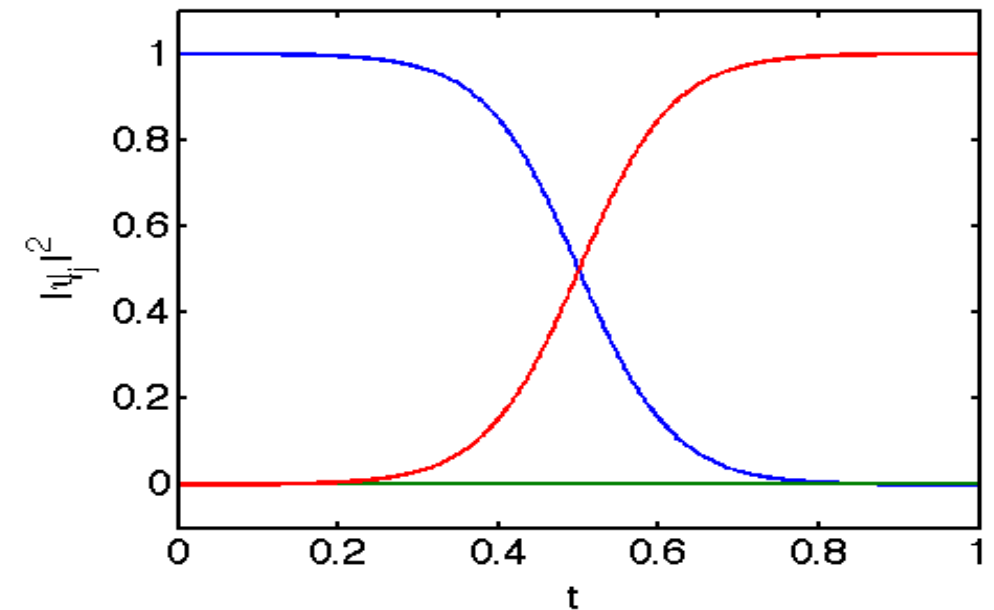
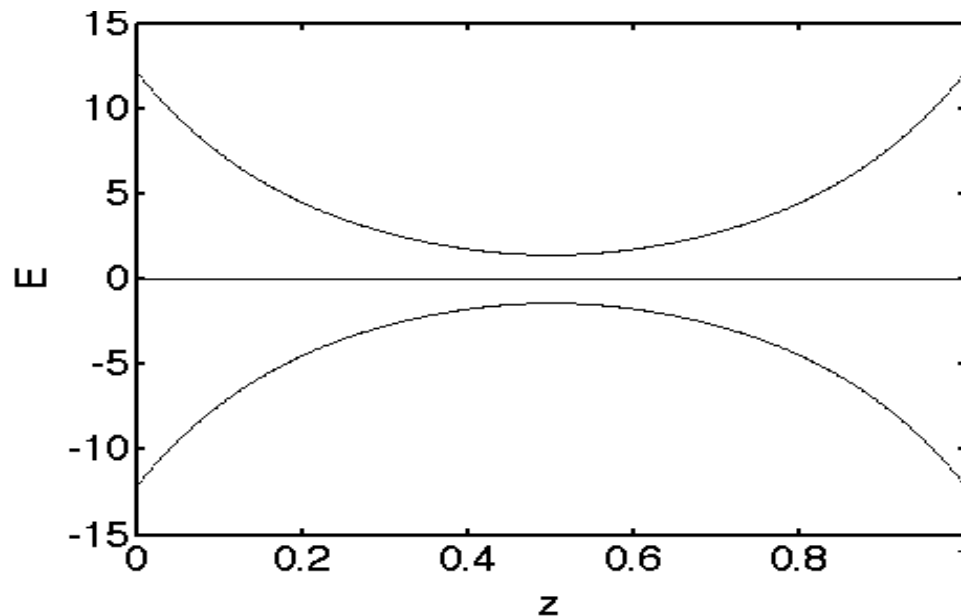
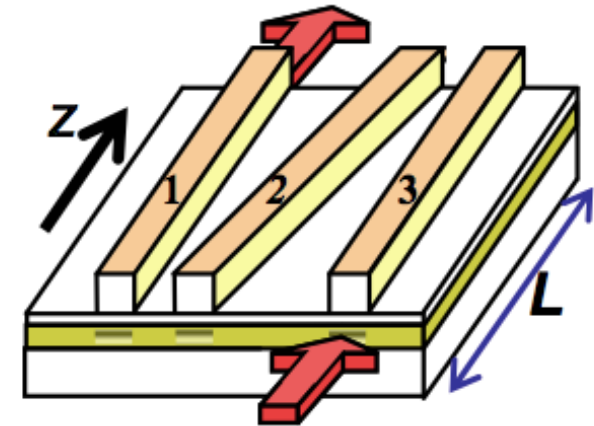
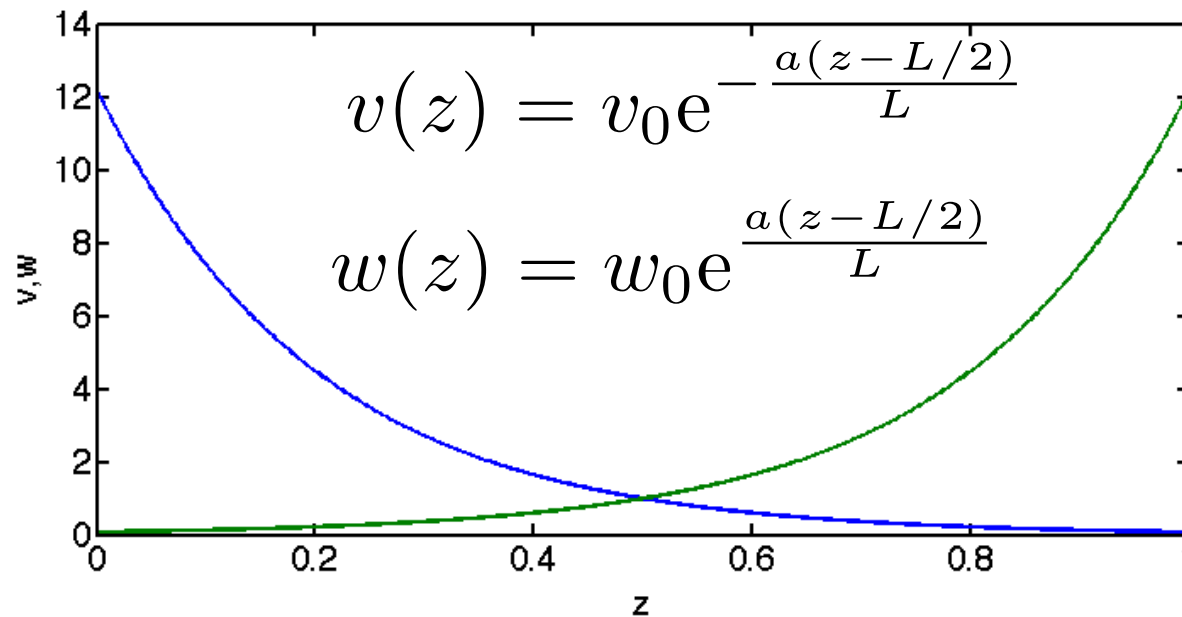
★ Eigenstates:

$$|0\rangle = \begin{pmatrix} \cos(\theta) \\ 0 \\ -\sin(\theta) \end{pmatrix}, \quad |\pm\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} \sin(\theta) \\ \pm 1 \\ \cos(\theta) \end{pmatrix}$$

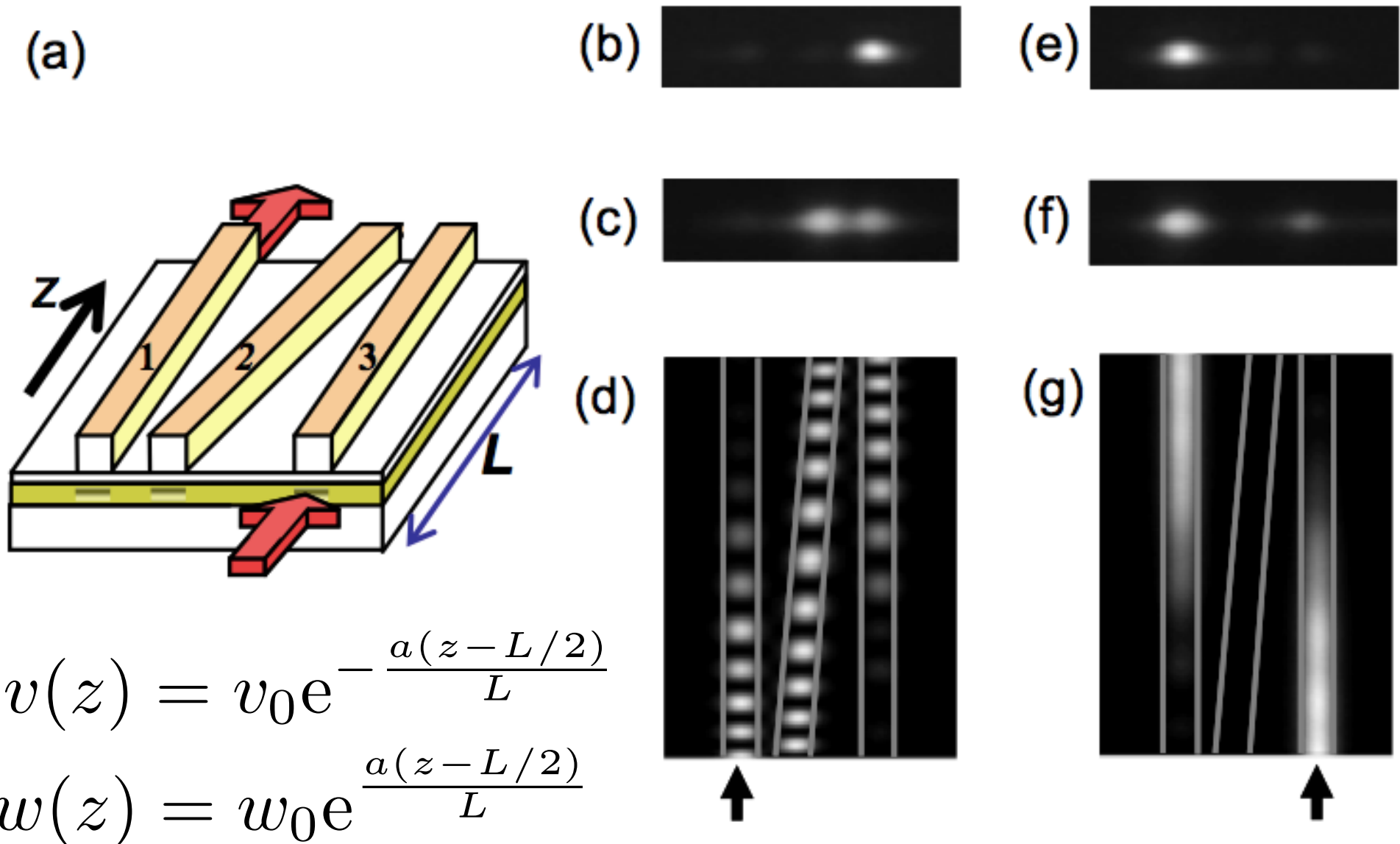
Adiabatic rotation from $|1\rangle$ to $|3\rangle$

$$\tan(\theta) = \frac{v}{w}$$

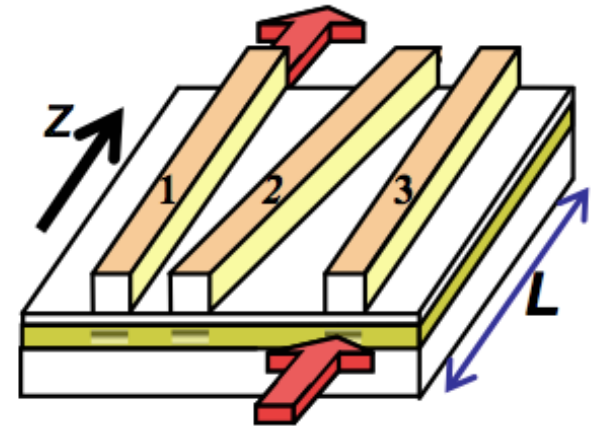
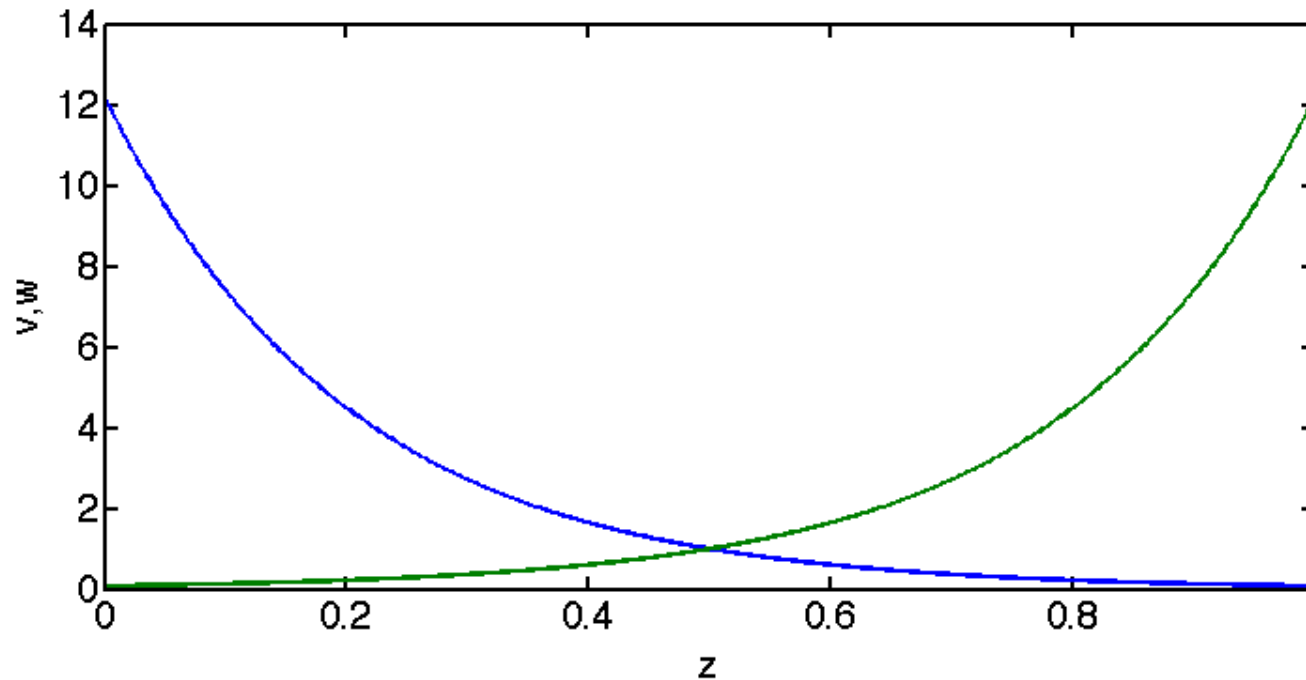
STIRAP type scheme in optical waveguides



STIRAP type scheme in optical waveguides



Additional decay/absorption



$$H = \begin{pmatrix} -i\gamma_1 & v(z) & 0 \\ v(z) & -i\gamma_2 & w(z) \\ 0 & w(z) & -i\gamma_3 \end{pmatrix}$$

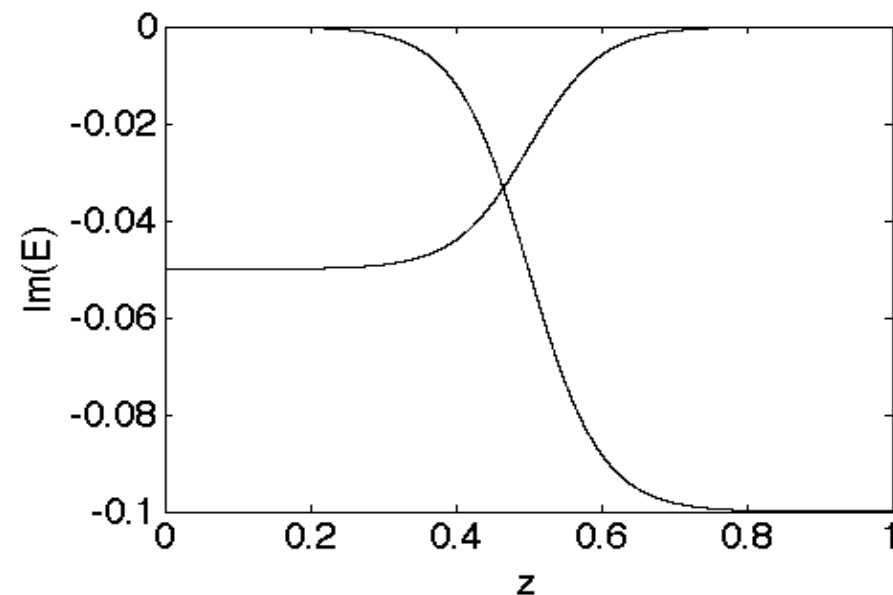
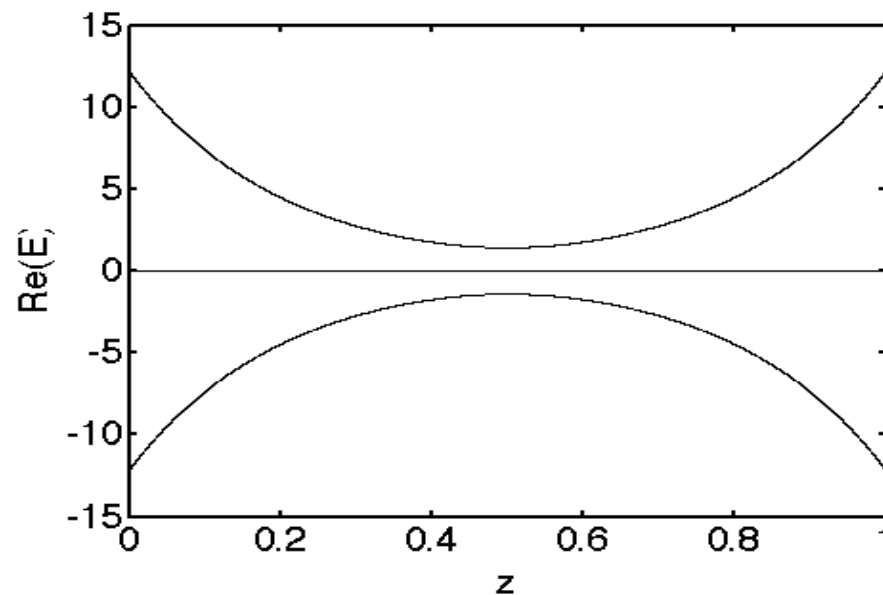
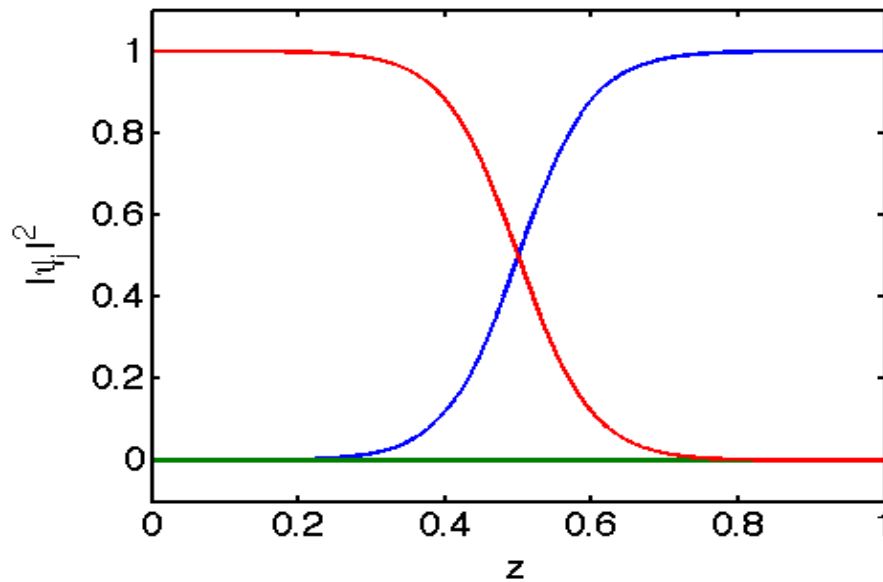
Absorption in
waveguides

Independent of γ_2

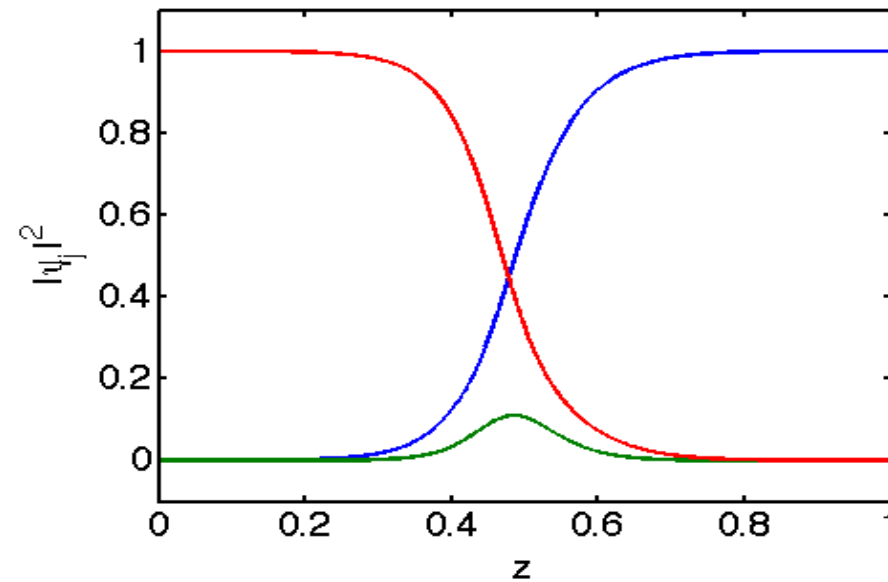
Decay/absorption in final state

$$\gamma_1 = 0.1, \quad \gamma_3 = 0$$

$$H = \begin{pmatrix} -i\gamma_1 & v(t) & 0 \\ v(t) & -i\gamma_2 & w(t) \\ 0 & w(t) & -i\gamma_3 \end{pmatrix}$$

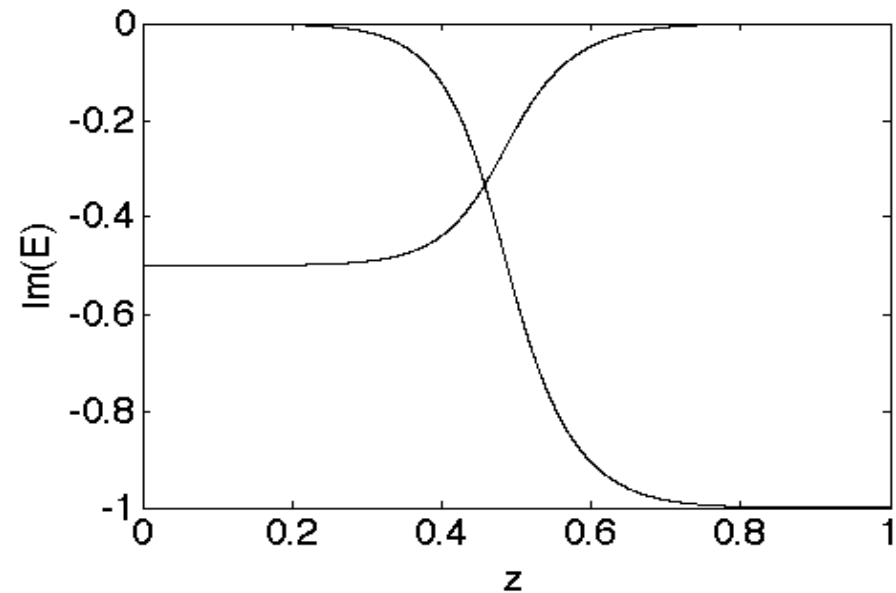
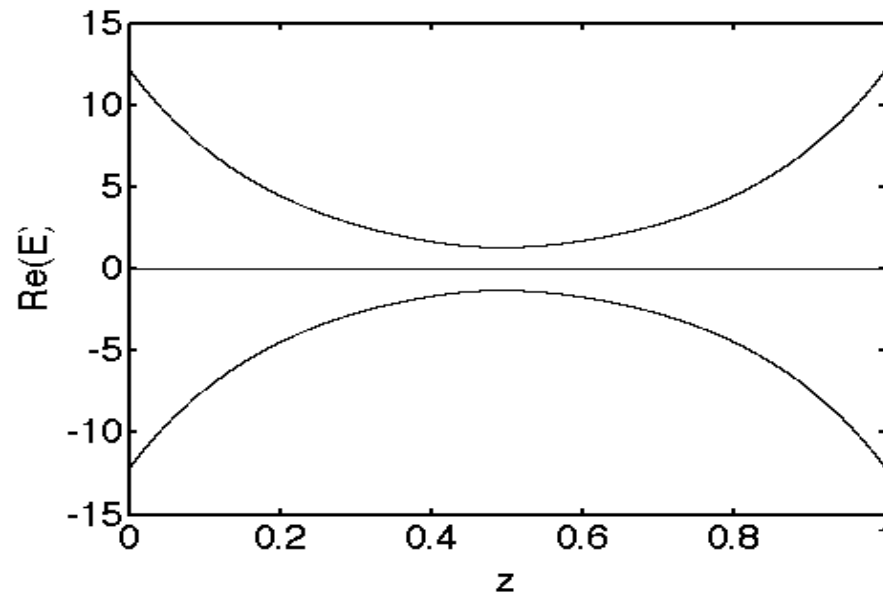


Decay/absorption in final state



$$\gamma_1 = 1, \quad \gamma_3 = 0$$

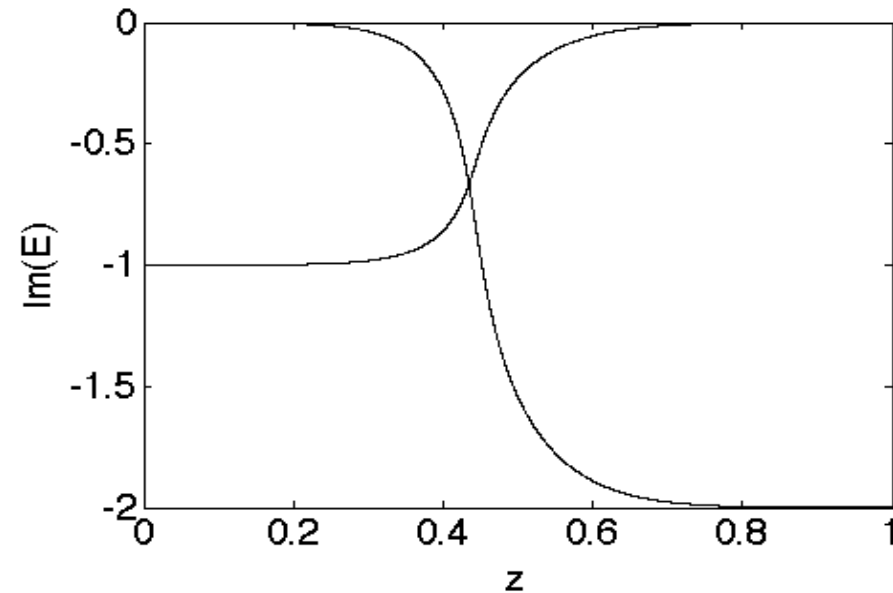
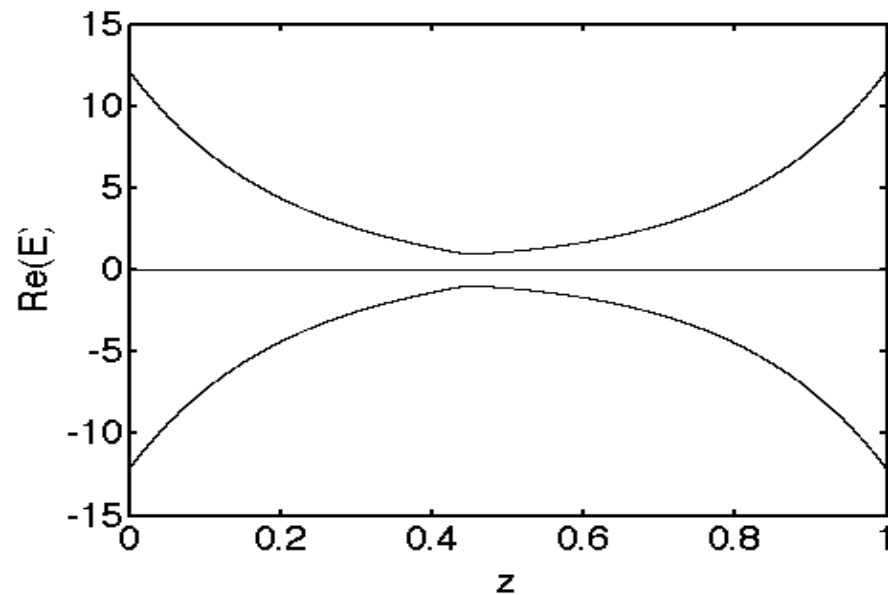
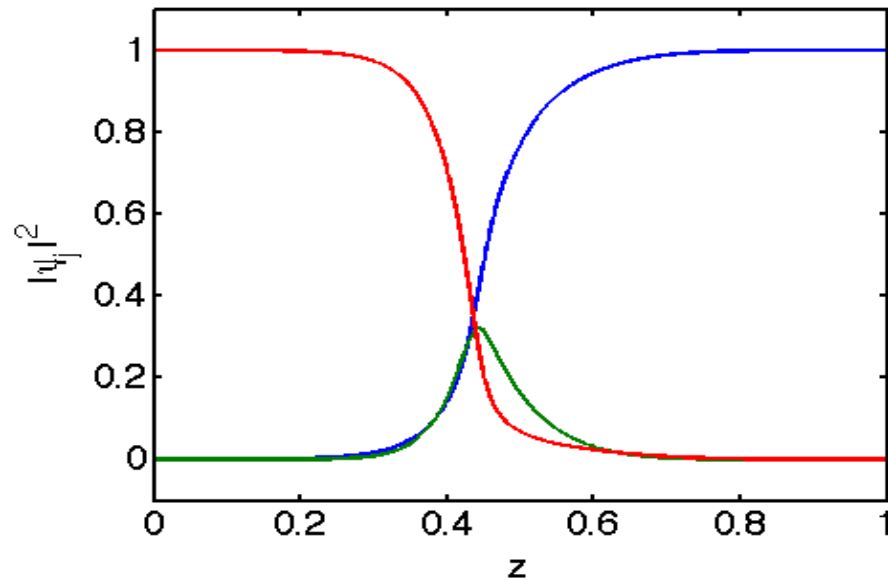
$$H = \begin{pmatrix} -i\gamma_1 & v(t) & 0 \\ v(t) & -i\gamma_2 & w(t) \\ 0 & w(t) & -i\gamma_3 \end{pmatrix}$$



Decay/absorption in final state

$$\gamma_1 = 2, \quad \gamma_3 = 0$$

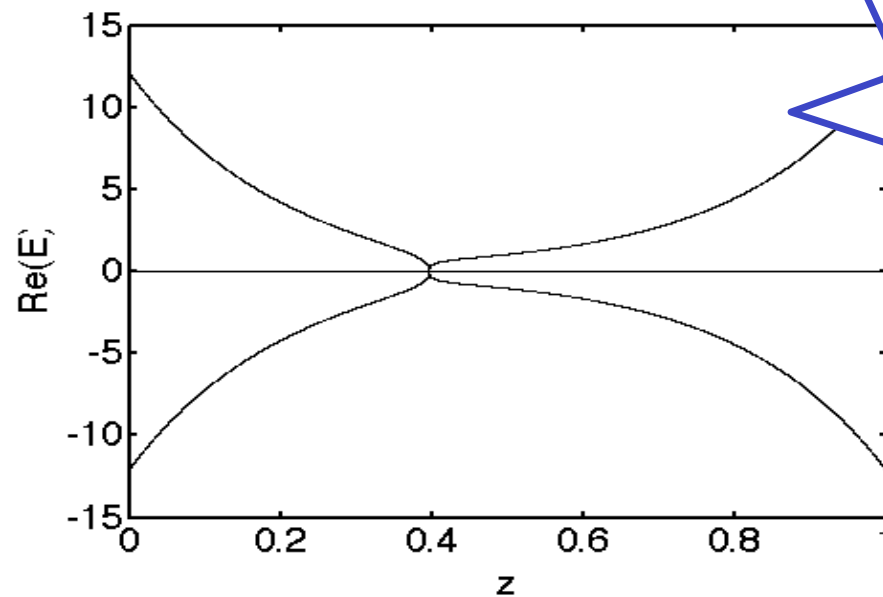
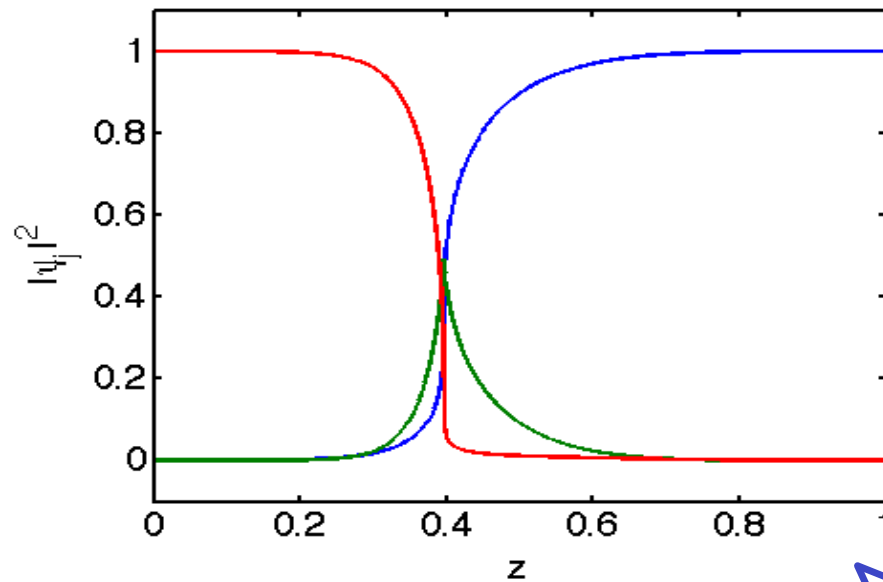
$$H = \begin{pmatrix} -i\gamma_1 & v(t) & 0 \\ v(t) & -i\gamma_2 & w(t) \\ 0 & w(t) & -i\gamma_3 \end{pmatrix}$$



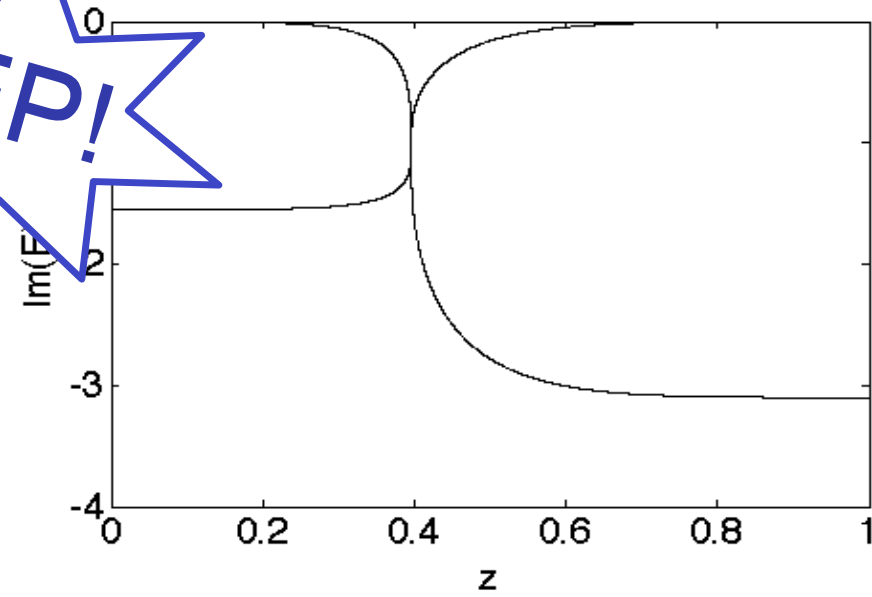
Decay/absorption in final state

$$\gamma_1 = 3.1, \quad \gamma_3 = 0$$

$$H = \begin{pmatrix} -i\gamma_1 & v(t) & 0 \\ v(t) & -i\gamma_2 & w(t) \\ 0 & w(t) & -i\gamma_3 \end{pmatrix}$$

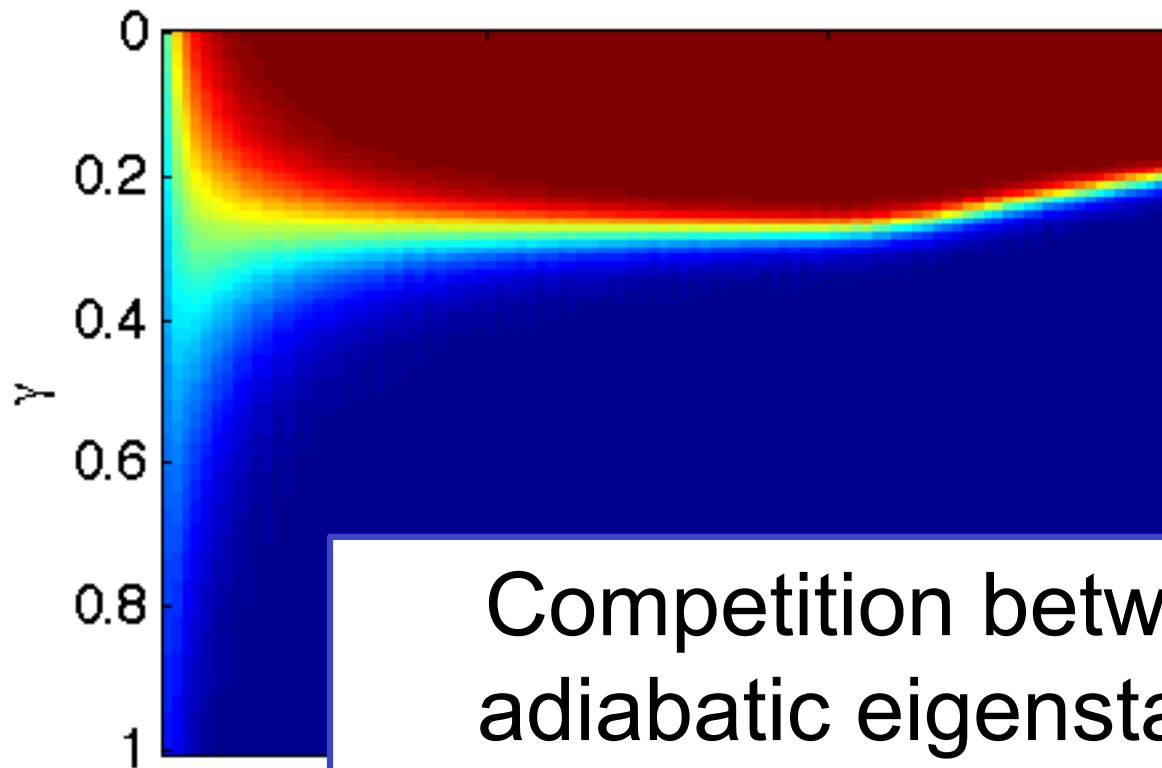


EP!



Transfer efficiency: absorption in final state

$$P = \frac{|\langle 1 | \psi(L) \rangle|^2}{\sum_{n=1}^3 |\langle n | \psi(L) \rangle|^2} \quad |\psi(0)\rangle = |\varphi^0(0)\rangle \approx -|3\rangle$$



Breakdown of
STIRAP long
before EP!

Competition between decay of
adiabatic eigenstates and non-
adiabatic transitions

Adiabatic quantum evolution

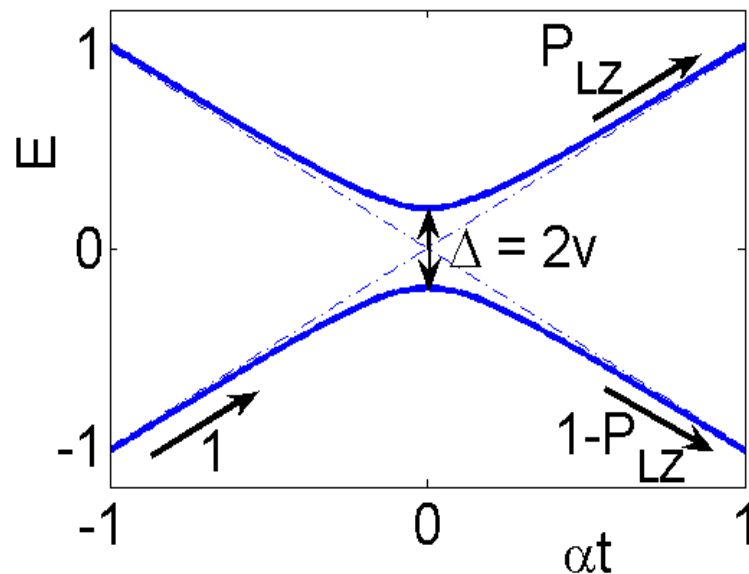
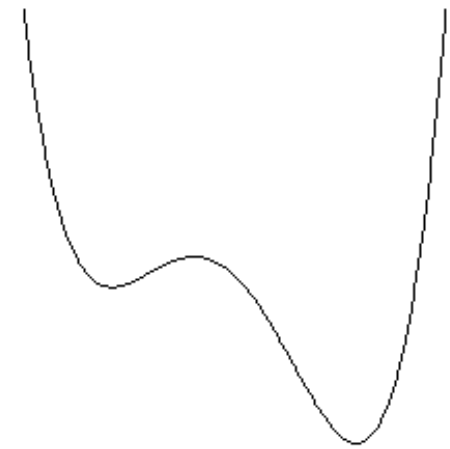
Adiabatic theorem: System prepared in eigenstate stays in corresponding instantaneous eigenstate for infinitely slow parameter variation

- ★ Adiabatic state transfer: Slow parameter variation such that initial and target states are connected via instantaneous states
- ★ Parameter variation in finite time: Small non-adiabatic corrections
- ★ Typical non-adiabatic corrections: Landau-Zener

Landau Zener model

$$i \frac{d}{dt} \begin{pmatrix} \psi_1 \\ \psi_2 \end{pmatrix} = \begin{pmatrix} \alpha t & v \\ v & -\alpha t \end{pmatrix}$$

$$t \in (-\infty, \infty)$$



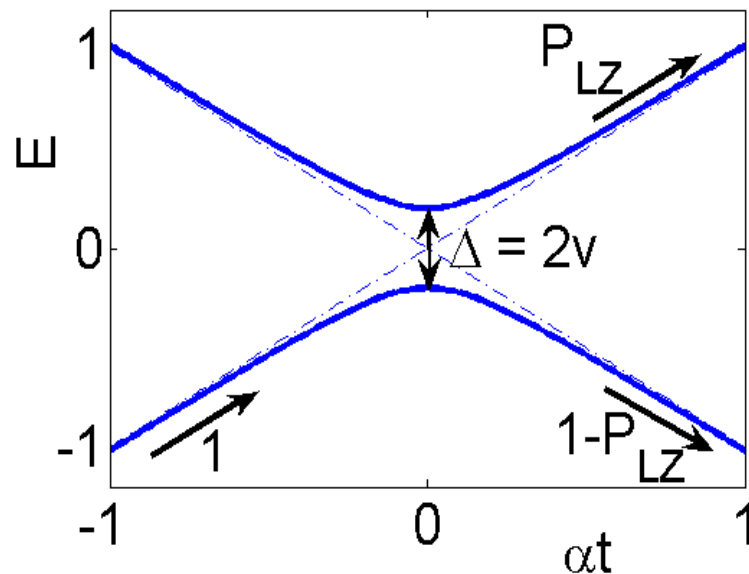
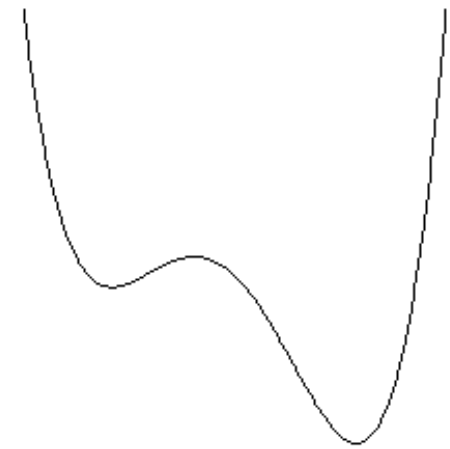
$$P_{LZ} = \frac{|\psi_1(t \rightarrow +\infty)|^2}{|\psi_1(t \rightarrow -\infty)|^2}$$

$$P_{LZ} = e^{-2 \text{Im} \int_0^{t_0} (E_+ - E_-) dt}$$

Landau Zener model

$$i \frac{d}{dt} \begin{pmatrix} \psi_1 \\ \psi_2 \end{pmatrix} = \begin{pmatrix} \alpha t & v \\ v & -\alpha t \end{pmatrix}$$

$$t \in (-\infty, \infty)$$



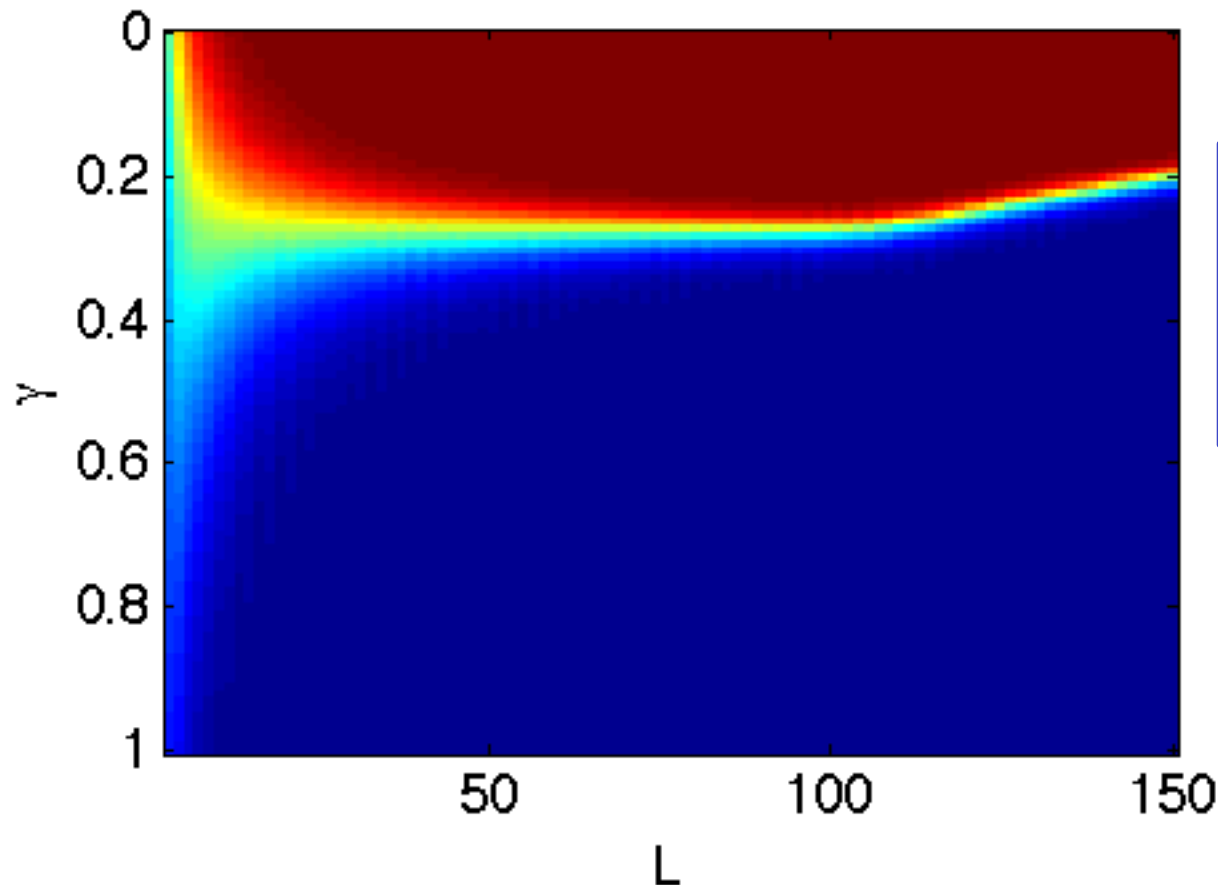
$$P_{LZ} = \frac{|\psi_1(t \rightarrow +\infty)|^2}{|\psi_1(t \rightarrow -\infty)|^2}$$

$$P_{LZ} = e^{-\pi v^2/|\alpha|}$$

$$\text{Adiabaticity: } P_{LZ} = 0 \text{ or } 1 \quad \alpha \rightarrow 0$$

Transfer efficiency: absorption in final state

$$P = \frac{|\langle 1 | \psi(L) \rangle|^2}{\sum_{n=1}^3 |\langle n | \psi(L) \rangle|^2} \quad |\psi(0)\rangle = |\varphi^0(0)\rangle \approx -|3\rangle$$



Breakdown of
STIRAP long
before EP!

Adiabatic transfer boundary

$$P = \frac{|\langle 1 | \psi(L) \rangle|^2}{\sum_{n=1}^3 |\langle n | \psi(L) \rangle|^2} \quad |\psi(0)\rangle = |\varphi^0(0)\rangle \approx -|3\rangle$$

★ Adiabatic and non-adiabatic contribution:

$$|\psi(z)\rangle = \psi_{ad}(z)|\varphi^0(z)\rangle + \sum_{j=\pm} \psi_{nonad}(z)|\varphi^j(z)\rangle$$

$$P = \frac{|\psi_{ad}(L)|^2}{|\psi_{ad}(L)|^2 + 2|\psi_{nonad}(L)|^2}$$

★ Estimate (non-)adiabatic contributions!

Adiabatic transfer boundary

$$|\psi(z)\rangle = \psi_{ad}(z)|\varphi^0(z)\rangle + \sum_{j=\pm} \psi_{nonad}(z)|\varphi^j(z)\rangle$$

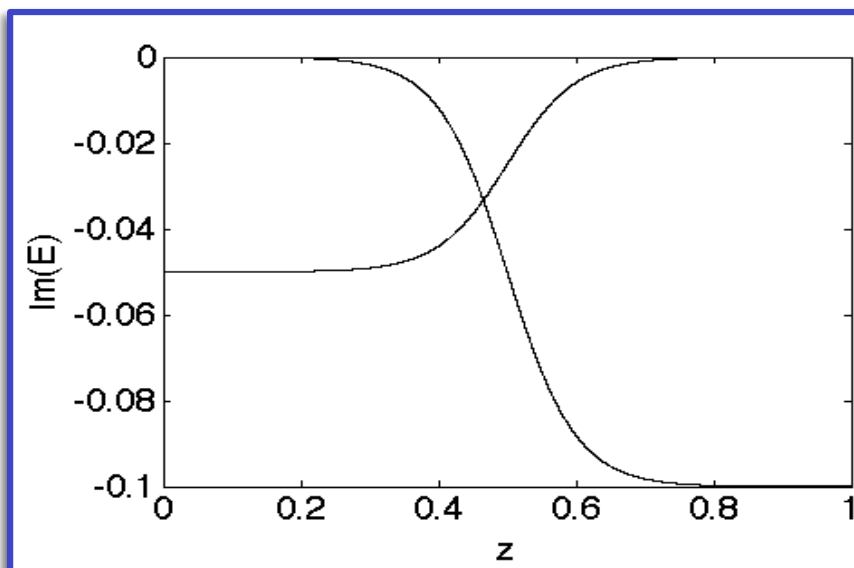
★ Transition happens around $z = L/2$

★ $|\varphi^\pm(z)\rangle$ almost stable after transition:

$$|\psi_{nonad}(L)| \approx \sqrt{P_{nonad}}$$

★ $|\varphi^0(z)\rangle$ almost stable before transition:

$$|\psi_{ad}(L)| = \sqrt{1 - 2P_{nonad}} \exp\left(\int_0^L \text{Im } E_0 dz\right) \\ \approx \sqrt{1 - 2P_{nonad}} e^{-\gamma L/2}$$



Adiabatic transfer boundary

$$P = \frac{|\psi_{ad}(L)|^2}{|\psi_{ad}(L)|^2 + 2|\psi_{nonad}(L)|^2}$$

★ $|\psi_{nonad}(L)| \approx \sqrt{P_{nonad}}$

★ $|\psi_{ad}(L)| \approx \sqrt{1 - 2P_{nonad}} e^{-\gamma L/2}$

$$P \approx \frac{1}{2} \Rightarrow \gamma \approx \ln \left(\frac{1}{2P_{nonad}} - 1 \right) / L$$

★ Landau-Zener type approximation for $\gamma \neq 0$

$$P_{nonad} \approx \exp \left(-\frac{2}{a\sqrt{\pi}} \Gamma^2 \left(\frac{3}{4} \right) L \right)$$

Adiabatic transfer boundary

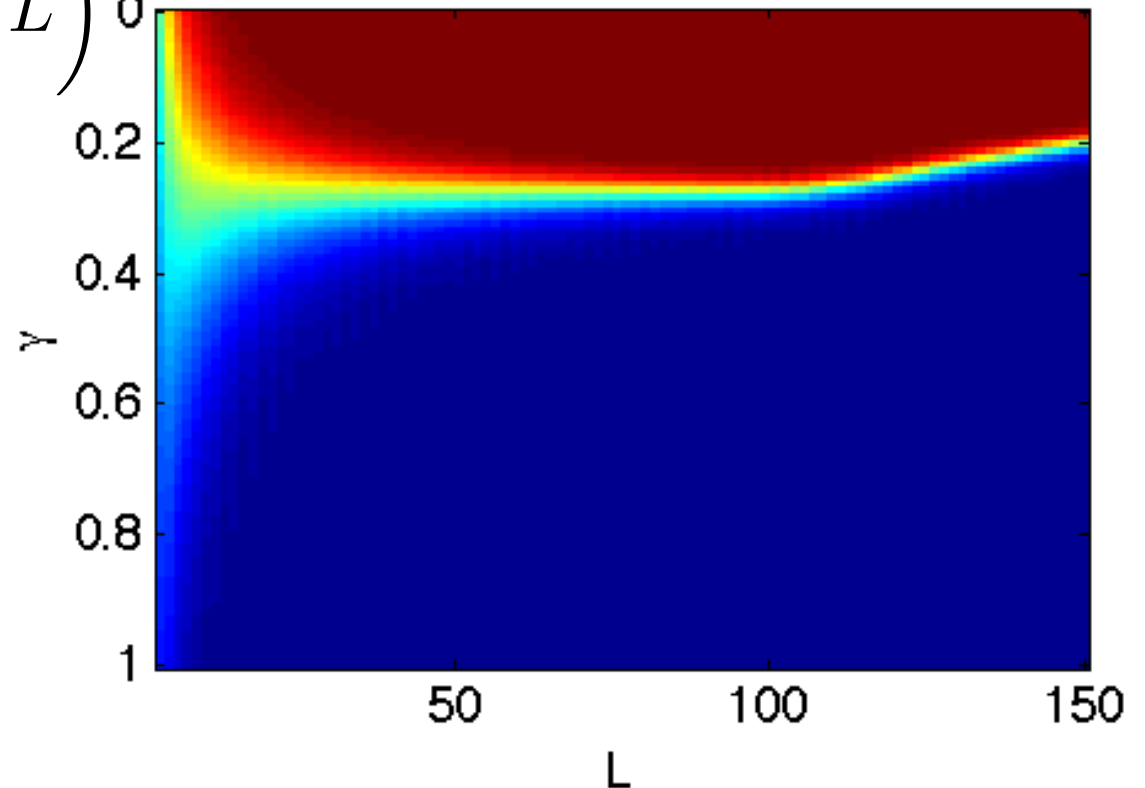
$$\gamma \approx \ln \left(\frac{1}{2P_{nonad}} - 1 \right) / L$$

$$P_{nonad} \approx \exp \left(-\frac{2}{a\sqrt{\pi}} \Gamma^2 \left(\frac{3}{4} \right) L \right)$$

$$\gamma_{cr}^{LZ} \approx \frac{2}{a\sqrt{\pi}} \Gamma^2 \left(\frac{3}{4} \right) - \frac{\ln(2)}{L}$$

★ Deviation due to LZ approximation and $\gamma \neq 0$

★ Use numerically obtained P_{nonad}



Adiabatic transfer boundary

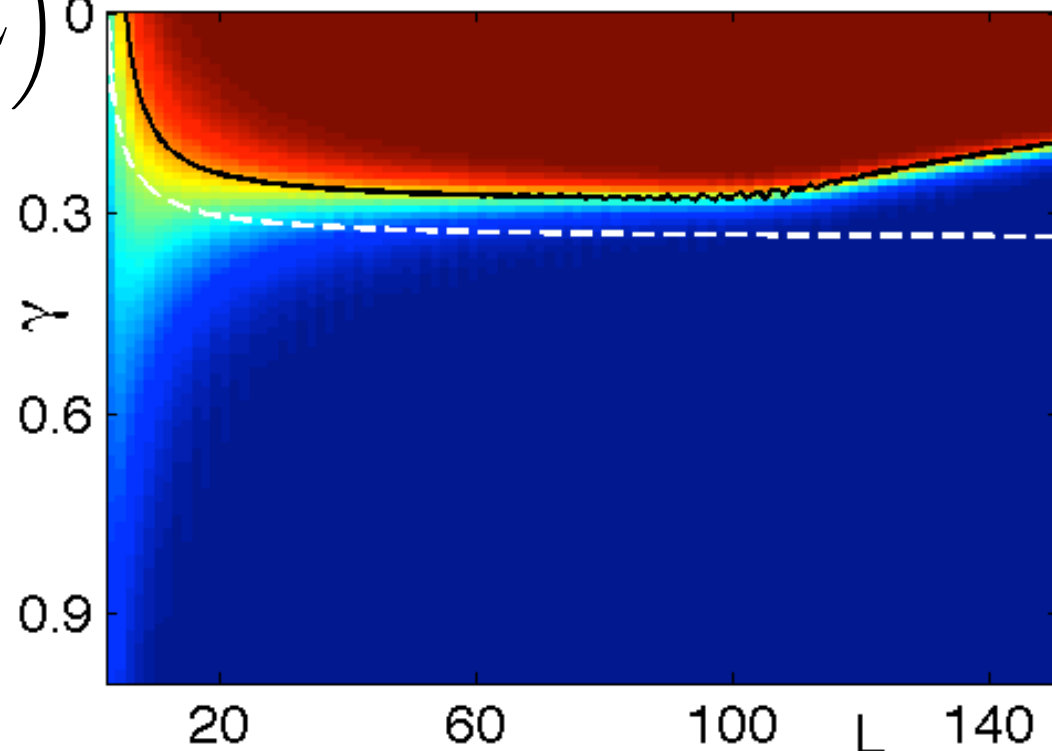
$$\gamma \approx \ln \left(\frac{1}{2P_{nonad}} - 1 \right) / L$$

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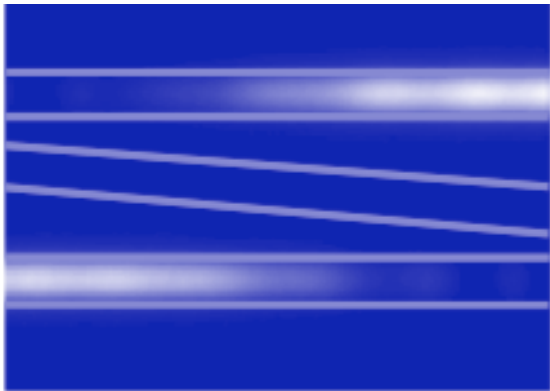
★ Deviation due to LZ approximation and $\gamma \neq 0$

★ Use numerically obtained P_{nonad}



Summary

- ★ Robust population transfer via adiabatic parameter variation



- ★ Example: STIRAP in optical waveguide structures

- ★ Even small absorption can destroy adiabaticity due to competition with non-adiabatic corrections

arxiv:1207.5235

Summary

- ★ Robust population transfer via adiabatic parameter variation

A Example: CTIRAD is optimal

Thank you for your attention!

- ★ Even small absorption can destroy adiabaticity due to competition with non-adiabatic corrections

arxiv:1207.5235